

LMDZ

Dynamics/physics organization,
Grids,
Time stepping,
Dissipation...

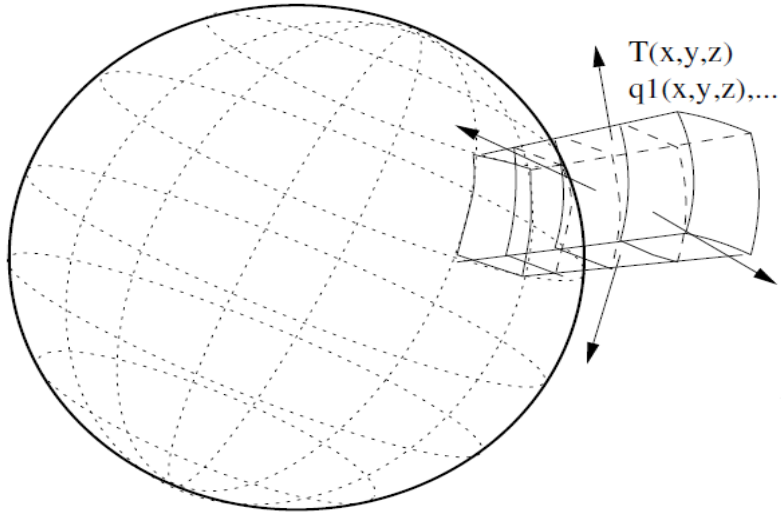
LMDZ courses, January 12, 2023

Overview of course topics

- **Grids:**
 - Horizontal grids in the physics & dynamics
 - Vertical discretization
- **Time marching:**
 - Generalities about time marching schemes
 - What is used in LMDZ
 - Longitudinal polar filter
- **Lateral diffusion and sponge layer:**
 - Energy cascade
 - Illustrative example of diffusion
 - Sponge layer near model top

Grids in LMDZ

Dynamics



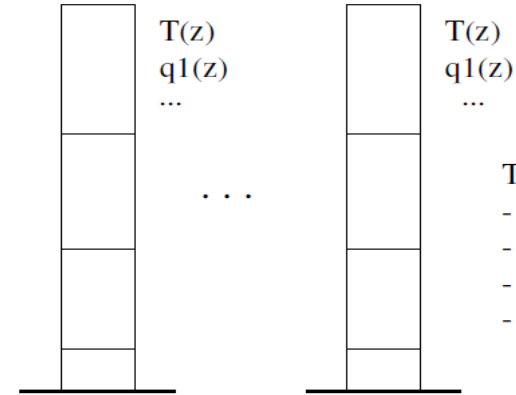
Dynamical tendencies



Physical fields



Physics



Tendencies due to:
- radiative transfer
- condensation
- subgrid dynamics
- ...

Separation between physics and dynamics:

- “**dynamics**”: solving the GFD equations on the sphere; usually with the assumption of a hydrostatic balance and thin layer approximation. Valid for all terrestrial planets.
- “**physics**”: (planet-specific) local processes, local to individual atmospheric columns.

Horizontal grids in LMDZ

Grid dimensions specified when compiling LMDZ:

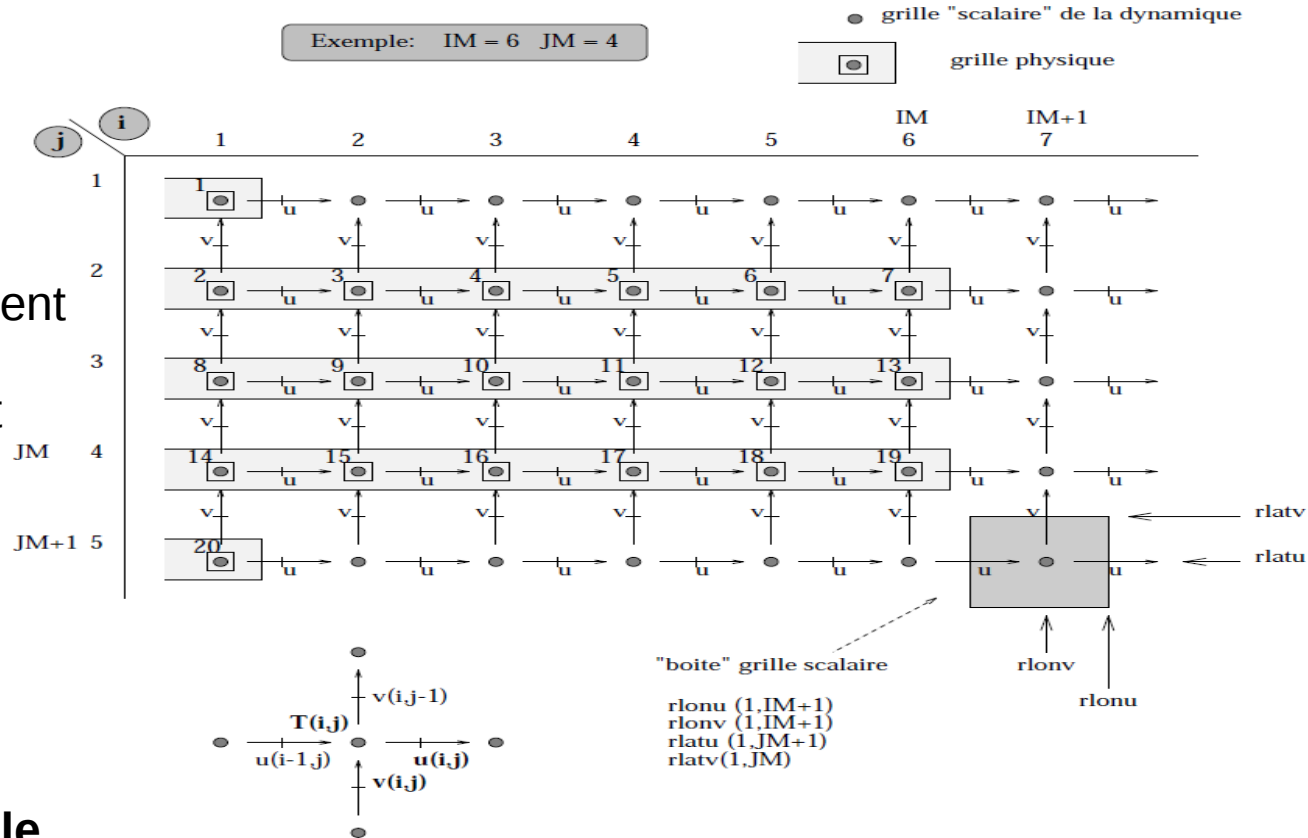
makeImdz[_fcm] -d iimxjjmxxllm ...

In the dynamics:

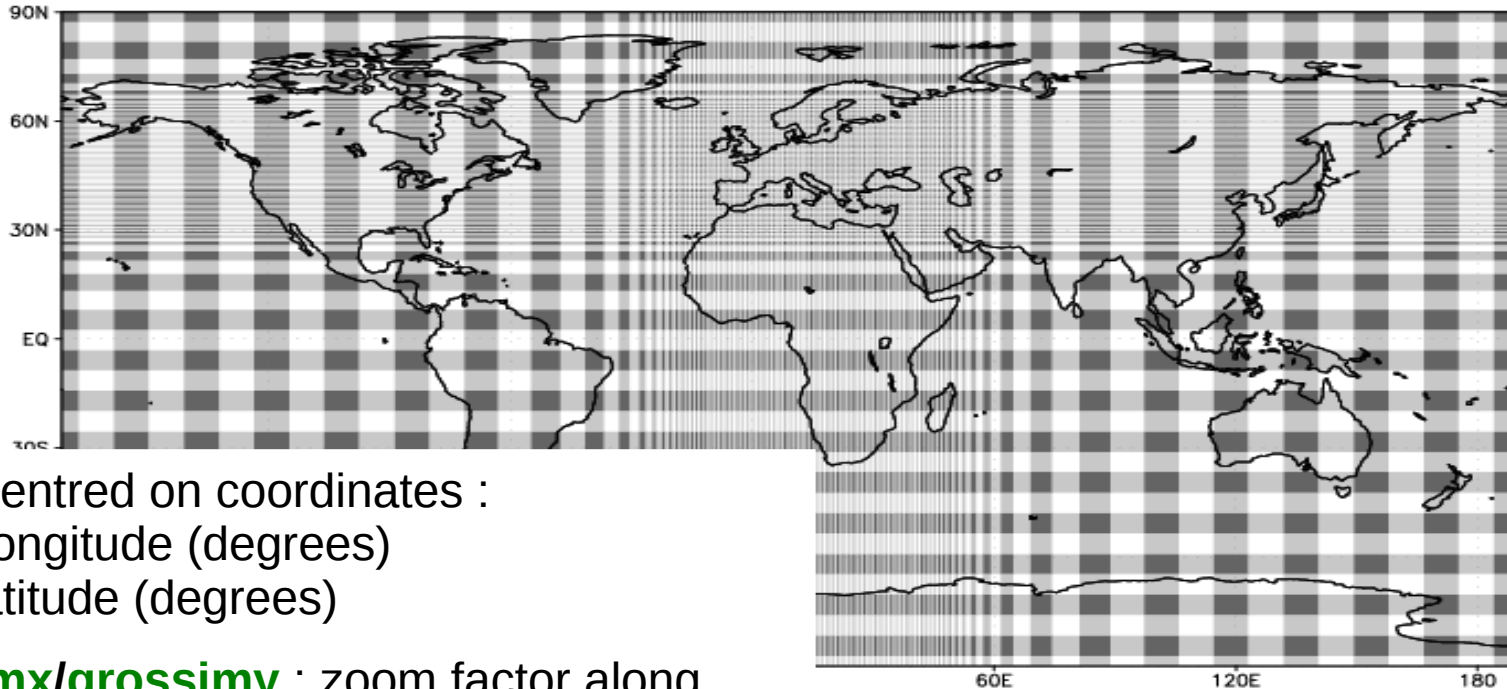
- Staggered grids: **u**, **v** and **scalars** (temperature, tracers) are on different meshes
- Global lonxlat grids with redundant grid points
 - at the poles
 - in longitude

In the physics:

- Collocated variables
- No global lonxlat horizontal grid, columns are labelled **using a single index** (from North Pole to South Pole)



LMDZ, Z for Zoom



Zoom centred on coordinates :

clon : longitude (degrees)

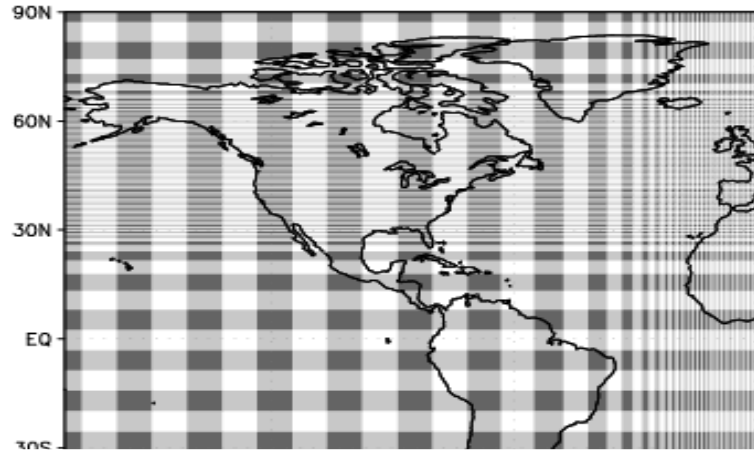
clat : latitude (degrees)

grossimx/grossimy : zoom factor along
x/y directions (i.e. lon/lat)

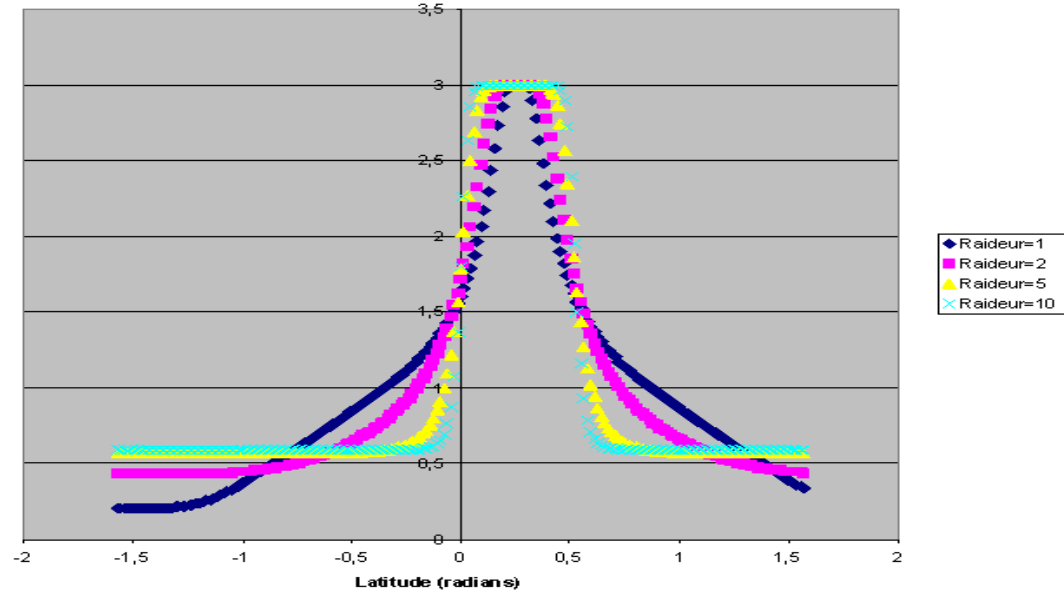
Computed as the ratio of the smallest mesh (i.e. in the zoom), compared to the mesh size for a global regular grid with the same total number of points.

dzoomx/dzoomy : fraction of the grid containing the zoomed area: $dzoom \times 360^\circ$ by $dzoomy \times 180^\circ$

LMDZ, Z for Zoom



GROSSISSEMENT LOCAL EN LATITUDE ($\varphi_0=15^\circ\text{N}$, $\Gamma=3$, $\Delta=30^\circ$)



Zoom centred on coordinates :

clon : longitude (degrees)

clat : latitude (degrees)

grossimx/grossimy : zoom factor along
x/y directions (i.e. lon/lat)

Computed as the ratio of the smallest mesh (i.e. in the zoom), compared to the mesh size for a global regular grid with the same total number of points.

taux/tauy : steepness of the transition between inner zoom and outer zoom meshes (typically one tries to avoid sharp transitions; $\tau \sim 3$ is a reasonable value)

Nudging in LMDZ

$$\frac{\partial u}{\partial t} = \frac{\partial u}{\partial t}_{GCM} + \frac{u_{analyse} - u}{\tau}$$
$$\frac{\partial v}{\partial t} = \frac{\partial v}{\partial t}_{GCM} + \frac{v_{analyse} - v}{\tau}$$

Nudging towards analyses or reanalyses with chosen time constants

(nudge = "guide" in French)

Example of nudging parameters:

ok_guide = y

guide_T = n , guide_p = n , guide_q = n

guide_u = y , guide_v = y

tau_min_u = 0.0208333 (days)

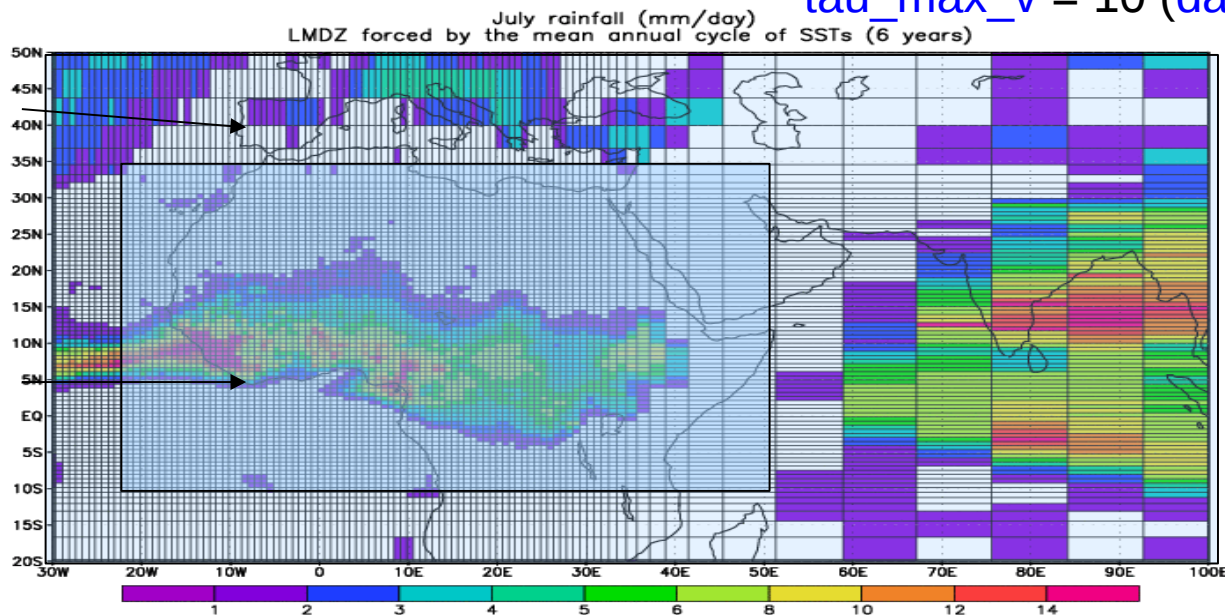
tau_max_u = 10 (days)

tau_min_v = 0.0208333 (days)

tau_max_v = 10 (days)

Strong nudging
($\tau=30\text{min}$)

Weak to moderate nudging
($\tau=10\text{ days}$)



Vertical discretization in LMDZ

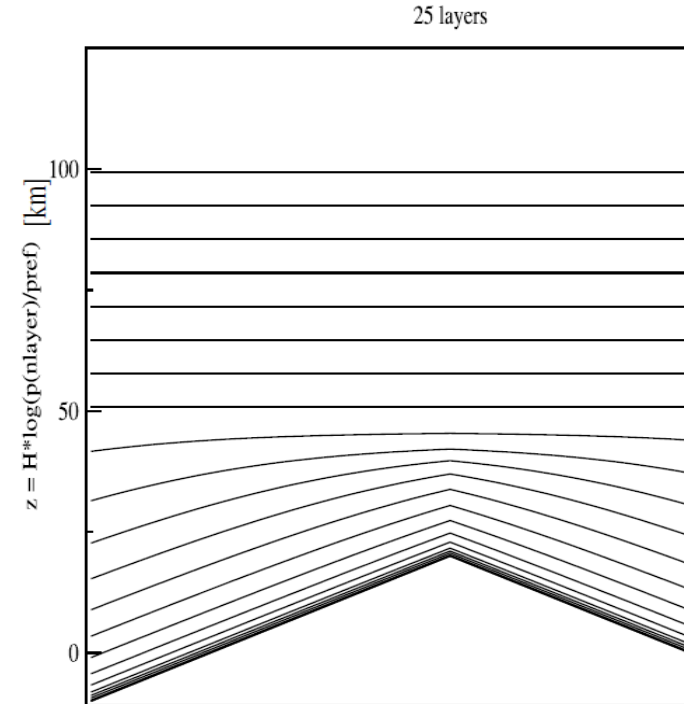
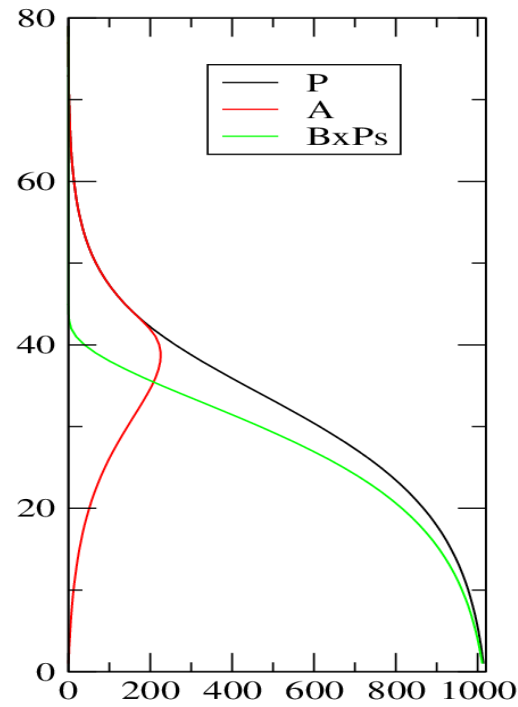
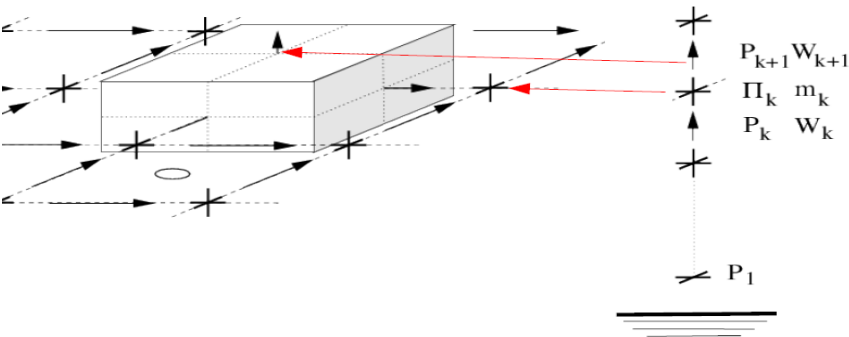
- Model levels are hybrid **sigma-pressure** levels:

$$P(\text{level}, \text{time}) = \text{ap}(\text{level}) + \text{bp}(\text{level}) \cdot P_s(\text{time})$$

hybrid coordinates $\text{ap}(k)$ and $\text{bp}(k)$ are fixed for a given model run

Surface pressure $P_s(t)$ varies during the run

- Near the surface $\text{ap} \sim 0$
 $\Rightarrow \text{bp}(k) \sim P/P_s$
- At high altitudes, $\text{bp} \sim 0$



Vertical discretization in LMDZ

- Setting model levels via the def files (also a function of number of vertical levels) :
vert_sampling = strato_custom : customizable (via other parameters in .def file; see next slide) discretization for stratospheric extensions.

Multiple other possibilities from this “default”:

vert_sampling = strato : a default for stratospheric extensions

vert_sampling = sigma : automated generation of purely sigma levels

vert_sampling = param : load values from a “sigma.def” file

vert_sampling = tropo : a default for tropospheric simulations

vert_sampling = read : read ap() and bp() from file “hybrid.txt”

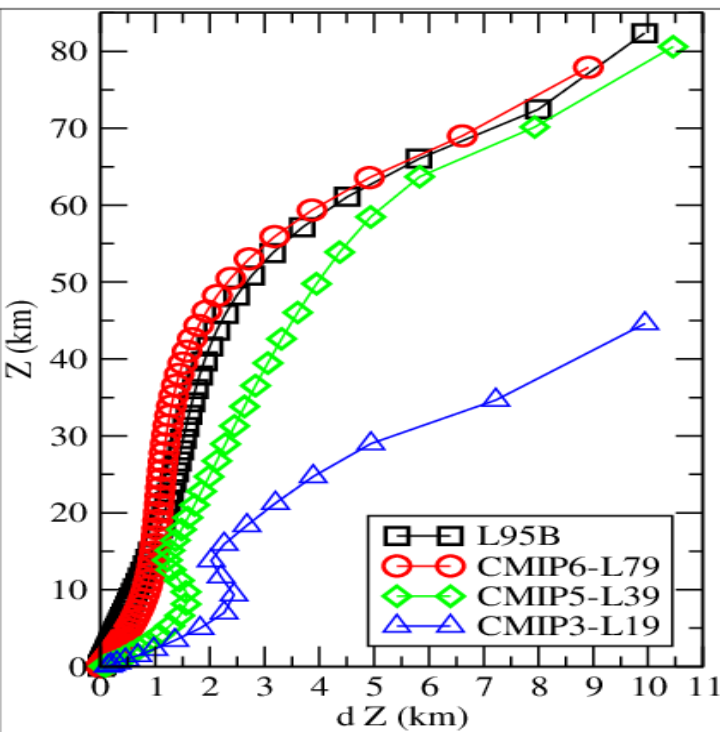
=> Typically you don't need to mess with the vertical discretization,

the default behaviour most likely matches your needs.

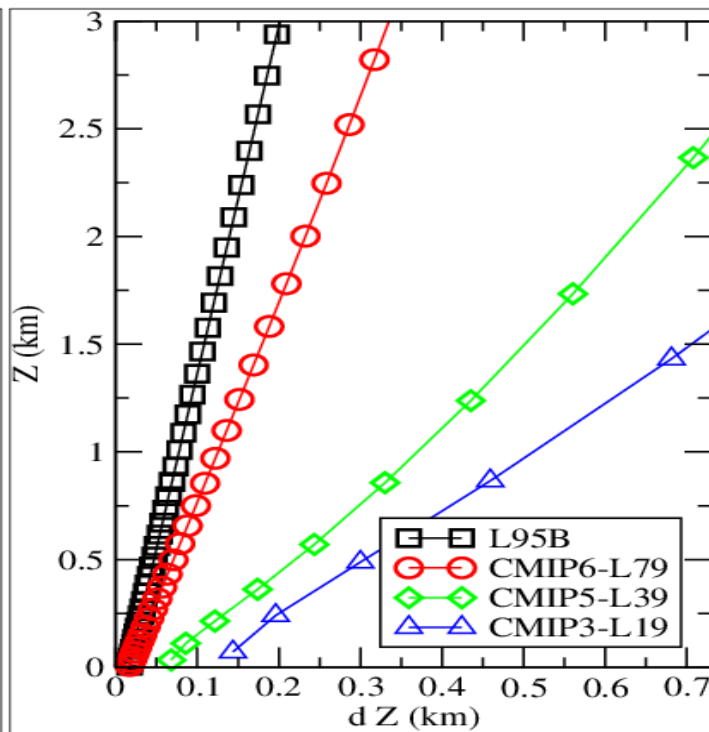
=> Check out routine dyn3d_common/**disvert.F90**

Vertical discretization in LMDZ

Illustration of typical altitudes and layer thickness of L19/L39/L79/L95 grids



Over the whole atmosphere
 $0 < z < \sim 80$ km



Near the surface
 $0 < z < 3$ km

**“Standard” CMIP6
L79 settings
(see DefLists/vert_L79.def)**

`vert_sampling=strato_custom`

`vert_scale=7.`
`vert_dzmin=0.017`
`vert_dzlow=1.`
`vert_z0low=8.7`
`vert_dzmid=2.`
`vert_z0mid=70.`
`vert_h_mid=20.`
`vert_dzhig=11.`
`vert_z0hig=75.`
`vert_h_hig=20.`

Questions ?

Time marching schemes

- **The big picture:** you want to solve

$$\frac{df(t)}{dt} = R(f, t)$$
$$f(t = 0) = f_0$$

from a known initial condition at time $t=0$ to time $t=...$

- So it is all about using a **time marching** scheme, built on **Taylor expansion** for evaluation of the time derivative, and choosing at which **time level** $t=n.dt$ the right hand side term $R[f(t),t]$ is to be evaluated

Time marching schemes

- **Explicit Euler** scheme (1st order in time):

$$\frac{df(t)}{dt} \simeq \frac{f_{n+1} - f_n}{\delta t}$$
$$R(f, t) \simeq R(f(t_n), t_n)$$

- **Implicit Euler** scheme (1st order in time):

$$\frac{df(t)}{dt} \simeq \frac{f_{n+1} - f_n}{\delta t}$$
$$R(f, t) \simeq R(f(t_{n+1}), t_{n+1})$$

- **Crank-Nicholson** scheme (2nd order in time):

$$\frac{df(t)}{dt} \simeq \frac{f_{n+1} - f_n}{\delta t}$$
$$R(f, t) \simeq \frac{R(f(t_{n+1}), t_{n+1}) + R(f(t_n), t_n)}{2}$$

Time marching schemes

- **Matsuno** scheme: a predictor-corrector (Euler explicit-Euler Implicit) scheme (1st order):

$$\begin{aligned}\frac{df(t)}{dt} &\simeq \frac{f_{n+1} - f_n}{\delta t} \\ f^p(t_{n+1}) &= f(t_n) + \delta t \cdot R(f(t_n), t_n) \\ R(f, t) &\simeq R(f^p(t_{n+1}), t_{n+1})\end{aligned}$$

- **Leapfrog** scheme: use encompassing time steps to evaluate the derivative (2nd order):

$$\begin{aligned}\frac{df(t)}{dt} &\simeq \frac{f_{n+1} - f_{n-1}}{2\delta t} \\ R(f, t) &\simeq R(f(t_n), t_n)\end{aligned}$$

Time marching schemes

- Illustrative example, on a decay equation (known solution!)

$$\frac{dq(t)}{dt} = -\frac{1}{\tau}q(t) \longrightarrow q(t) = q_0 e^{-\frac{t}{\tau}}$$

- Building **Euler explicit** (E) & **implicit** (I) schemes:

$$\begin{aligned} \frac{dq(t)}{dt} &= -\frac{1}{\tau}q(t) \\ \Rightarrow \frac{q^{n+1} - q^n}{\delta t} &\simeq -\frac{1}{\tau}q^n \quad (\text{E.E.}) \\ \Leftrightarrow q^{n+1} - q^n &= -\frac{\delta t}{\tau}q^n \\ \Leftrightarrow q^{n+1} &= \left[1 - \frac{\delta t}{\tau}\right] q^n \end{aligned}$$

$$\begin{aligned} \frac{dq(t)}{dt} &= -\frac{1}{\tau}q(t) \\ \Rightarrow \frac{q^{n+1} - q^n}{\delta t} &\simeq -\frac{1}{\tau}q^{n+1} \quad (\text{E.I.}) \\ \Leftrightarrow q^{n+1} - q^n &= -\frac{\delta t}{\tau}q^{n+1} \\ \Leftrightarrow \frac{\tau + \delta t}{\tau}q^{n+1} &= q^n \\ \Leftrightarrow q^{n+1} &= \frac{1}{1 + \frac{\delta t}{\tau}}q^n \end{aligned}$$

Time marching schemes

- Illustrative example, on a decay equation

$$\frac{dq(t)}{dt} = -\frac{1}{\tau}q(t) \longrightarrow q(t) = q_0 e^{-\frac{t}{\tau}}$$

- Resulting integration schemes:

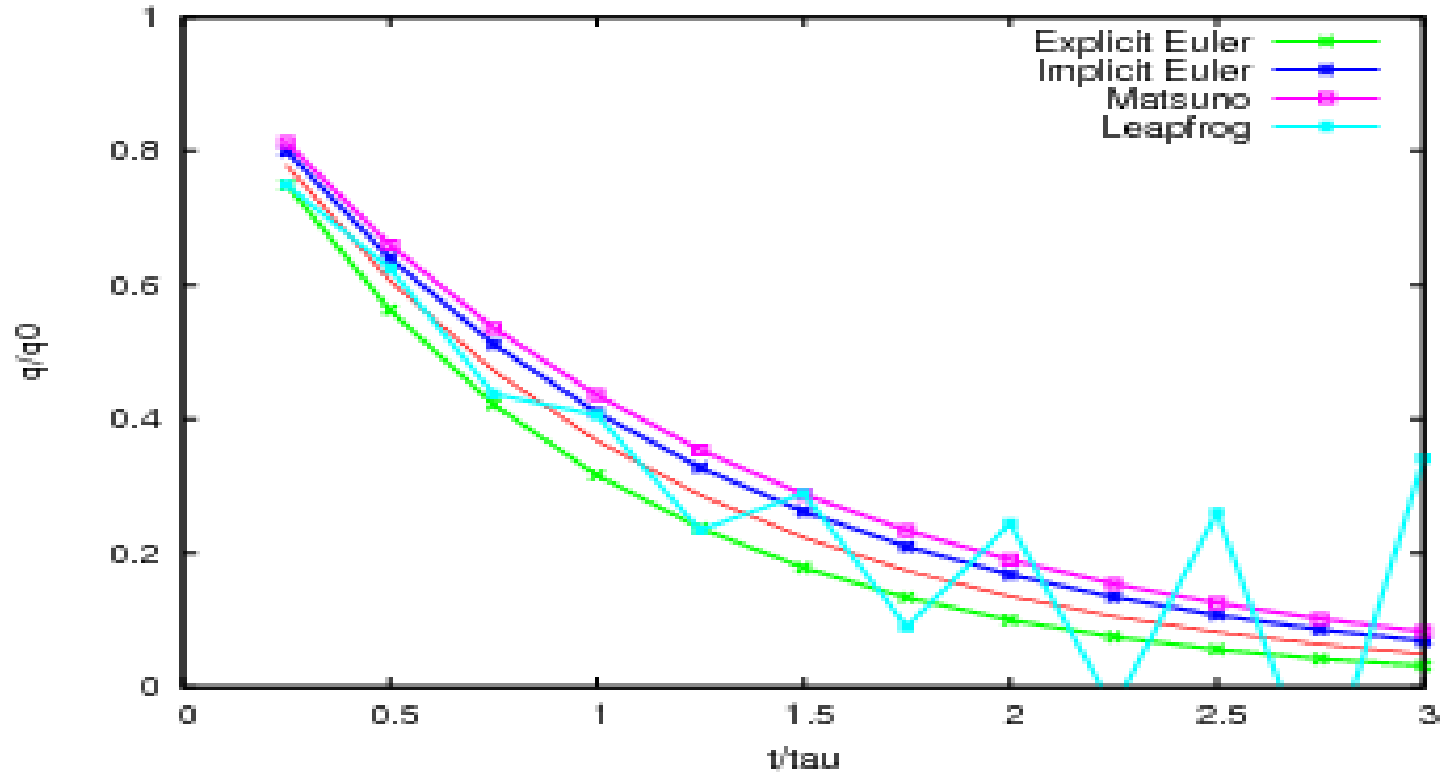
$$\text{EE : } q^{n+1} = \left[1 - \frac{\delta t}{\tau} \right] q^n$$

$$\text{EI : } q^{n+1} = \left[\frac{1}{1 + \delta t/\tau} \right] q^n$$

$$\text{CN : } q^{n+1} = \left[\frac{1 - \delta t/(2\tau)}{1 + \delta t/(2\tau)} \right] q^n$$

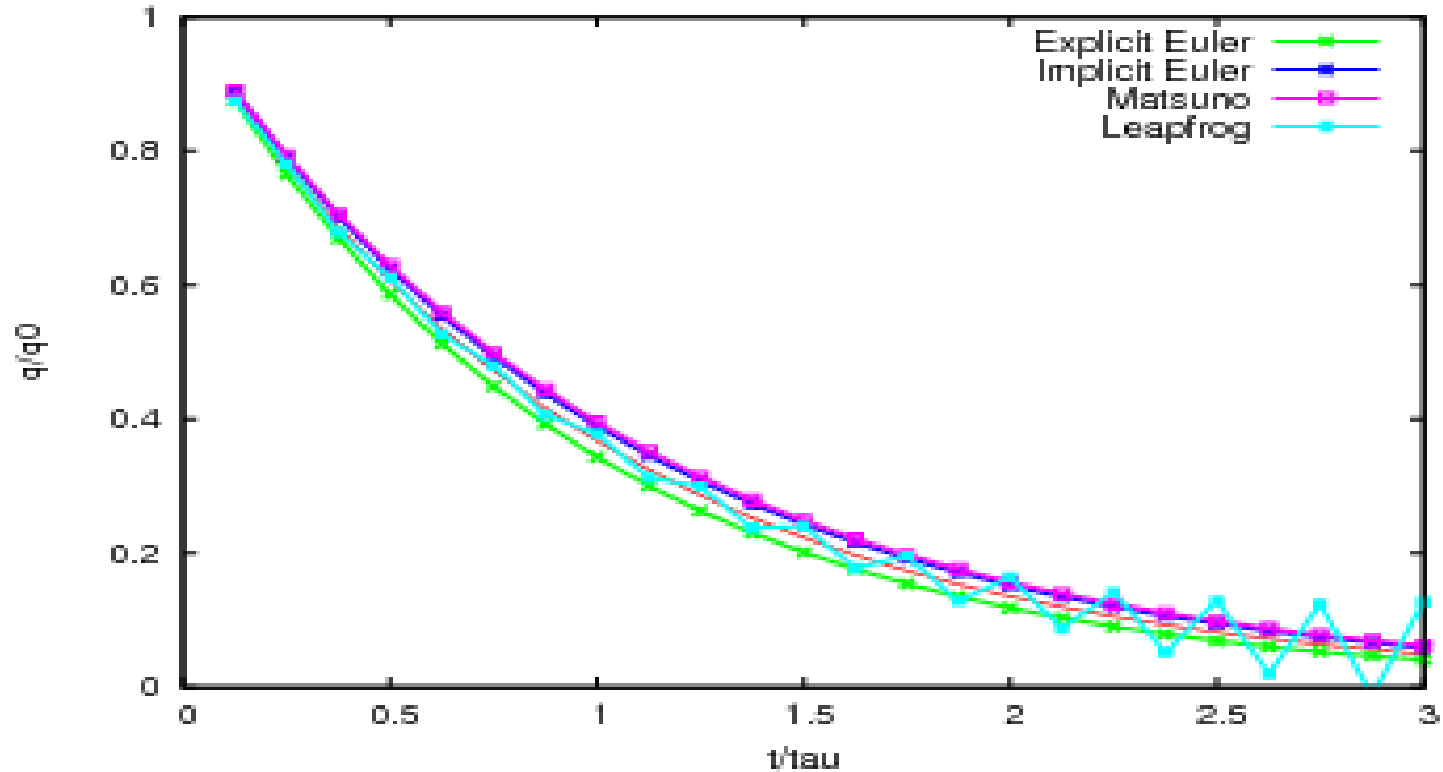
- Stability requirement (CFL) for EE : $dt/\tau < 2$

Time marching schemes



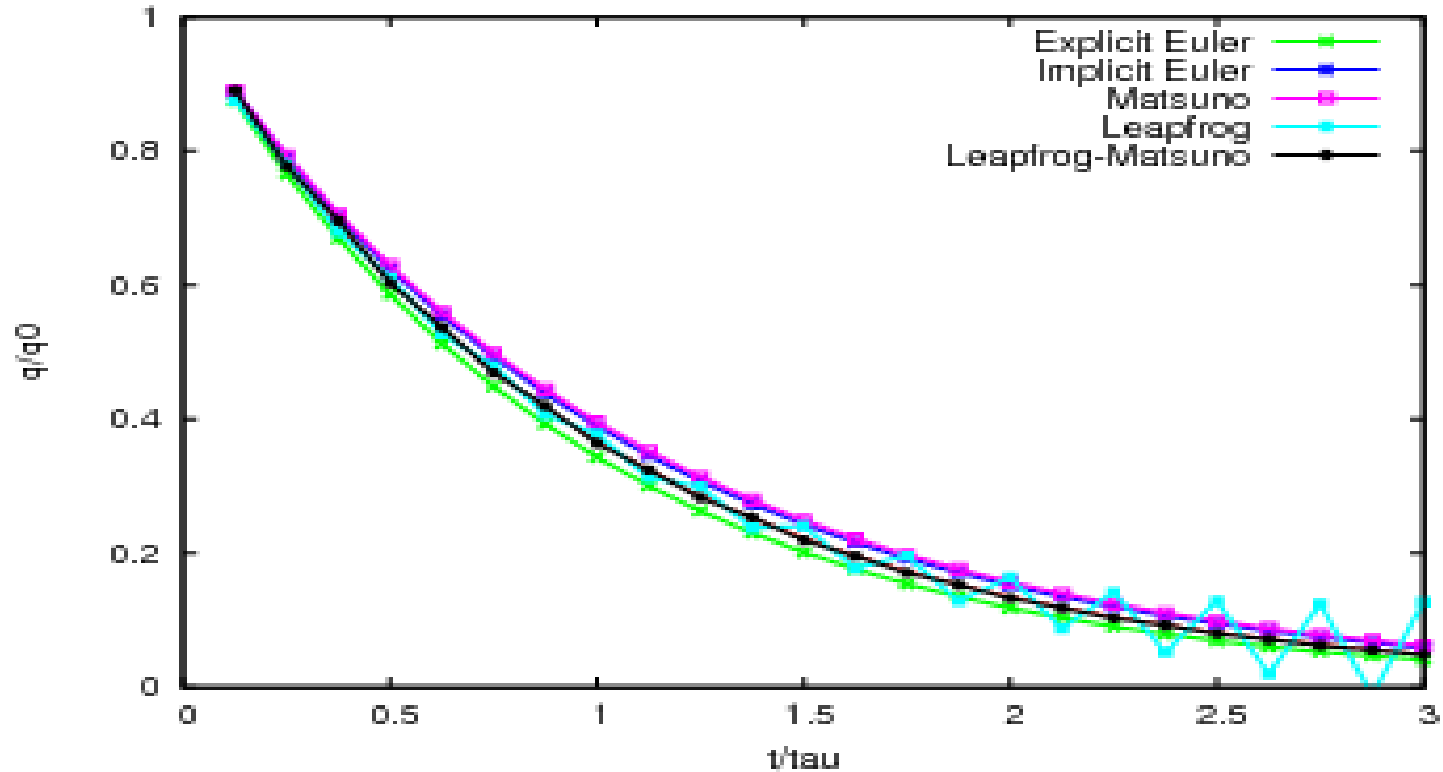
- 4 integration steps per unit of t/τ

Time marching schemes



- 8 integration steps per unit of t/τ

Time marching schemes



- 8 integration steps per unit of t/τ

Time marching in LMDZ

Time splitting between physics/dynamics/dissipation:

$$\frac{\partial \psi}{\partial t} = \textit{Dyn}(\psi) + \textit{Phy}(\psi) + \textit{Dissip}(\psi)$$

- **Dynamics** : Leapfrog-Matsuno scheme

Using `day_step` dynamical steps per day

Leapfrog steps with a Matsuno step every `iperiod` step

- **Physics** : Explicit Euler

Every `iphysiq` dynamical steps (multiple of `iperiod`)

- **Dissipation**: Explicit Euler

Every `dissip_period` dynamical steps (multiple of `iperiod`)

Side note about explicit or implicit time marching schemes

Even when solving linear spatio-temporal boundary-value problems, e.g.:

$$\frac{dA}{dt} = \kappa \frac{\partial^2 A}{\partial x^2}$$

The **explicit Euler** approach leads to a straightforward expression for grid point values (but with stability constraints) :

$$\frac{A_i^{k+1} - A_i^k}{\delta t} = \frac{\kappa}{h^2} [A_{i-1}^k - 2A_i^k + A_{i+1}^k]$$

Whereas the **implicit Euler** approach leads to a (tridiagonal) system of equations to solve:

$$\frac{A_i^{k+1} - A_i^k}{\delta t} = \frac{\kappa}{h^2} [A_{i-1}^{k+1} - 2A_i^{k+1} + A_{i+1}^{k+1}]$$

=> **requires more computations**, but may be necessary if time-stepping constraints require using large time steps.

Side note about tridiagonal system solving

When needing to solve a tridiagonal system of the form:

$T \cdot x = y$, T tridiagonal matrix, x & y vectors

Rather than invert T (costly!) to generate T^{-1} (dense matrix) and then compute $x = T^{-1} \cdot y$ (matrix-vector product)

Use the LU decomposition (Gaussian elimination) of T to split the problem into two very simple sub-problems:

- 1) $L \cdot U = T$, L and U are bidiagonal (lower/upper) matrices
- 2) Solve $L \cdot z = y$ for vector z (**forward substitution** step)
- 3) Solve $U \cdot x = z$ for vector x (**backward substitution** step)

Tracer advection in LMDZ

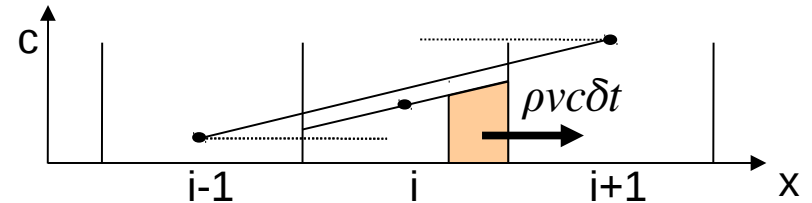
Use of the [Van Leer I scheme](#) (1977), a second order finite volume scheme with slope limiters (e.g. MUSCL, MINMOD) (Hourdin et Armengaud, 1999).

Guaranties of fundamental physical properties of transport :

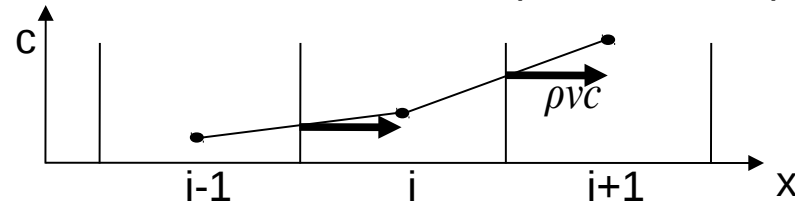
conservation of the total quantity,
positivity, monotony, non amplification
of extrema, weak numerical diffusion

- **CFL requirement**, for an advection velocity U_{max} :
 $U_{max} \cdot (dt/dx) = cte$, with $cte \sim 0(1)$

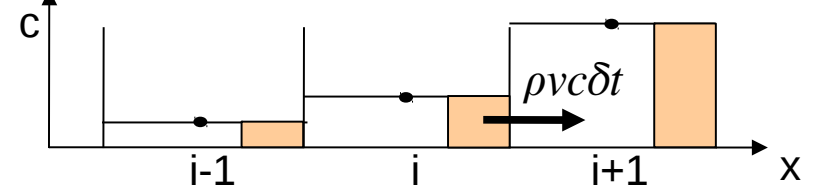
Scheme I by Van Leer (1977)



Centered finite differences (second order)



Upwind first order scheme (Godunov, 1952)



Tracer advection in LMDZ

- In practice: Tracer names and various properties are set in the `tracer.def` file. e.g.:

```
&version=1.0
```

```
&lmdz
```

```
default type=tracer phases=g hadv=10 vadv=10 parent=air
```

```
H2O phases=g hadv=14 vadv=14
```

```
H2O phases=l hadv=10 vadv=10
```

```
H2O phases=s hadv=10 vadv=10
```

```
Rn
```

```
Pb
```

- Scheme “10” : Van Leer scheme
- Scheme “14” : Dedicated modified scheme for water vapor
- Other (experimental) schemes are coded; see [dyn3d/advtrac.F90](#)

Tracer advection in LMDZ

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H2O phases=g hadv=14 vadv=14
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```
H2O phases=l hadv=10 vadv=10
```

```
H2O phases=s hadv=10 vadv=10
```

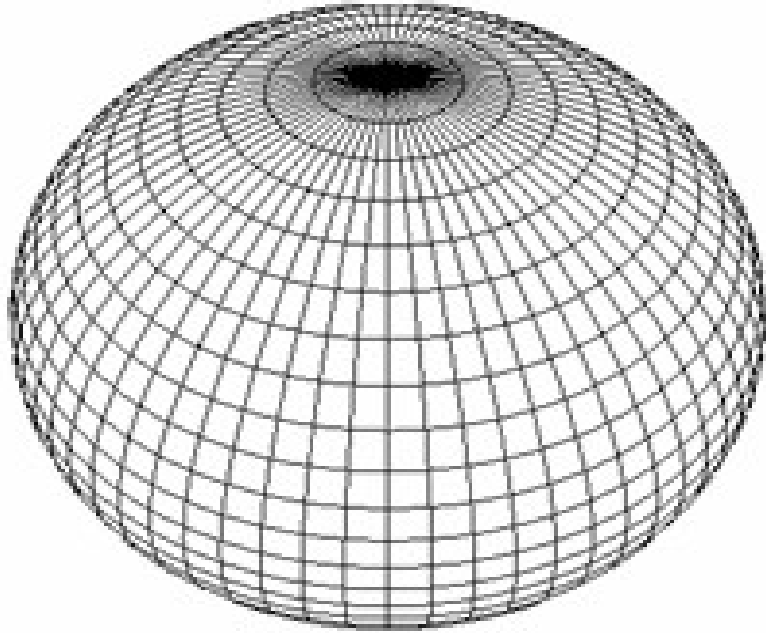
```
Rn
```

```
Pb
```

- phases: gas/liquid/solid for a given tracer
- parent: air or another tracer (for isotopes)+ possibility of “tagging”
see <https://lmdz-forge.lmd.jussieu.fr/mediawiki/LMDZPedia>

Questions ?

The longitudinal polar filter



- A lon-lat grid implies that the **meshes tighten dramatically** as the pole is approached.
- **CFL conditions** there would dictate using an extremely small time step for the time marching scheme.
- **Longitudinal (Fourier) filtering**, removing high spatial frequencies, is used to enforce that resolved features are at the level of those at $\sim 60^\circ$ (where longitudinal resolution is half of that at the equator).
- In addition near the poles there is some longitudinal grouping of meshes (applied to the divergence of air transport) by bunches of 2^{ngroup} (typically $\text{ngroup}=3$) which implies that the number of points along longitudes of the GCM must be a multiple of 2^{ngroup} !

Energy spectra and lateral dissipation

- Observations (Nastrom & Gage 1985, Lindborg 1999) collected over length scales from a few to thousands of km display a **characteristic energy cascade** (from Skamarock, 2004).

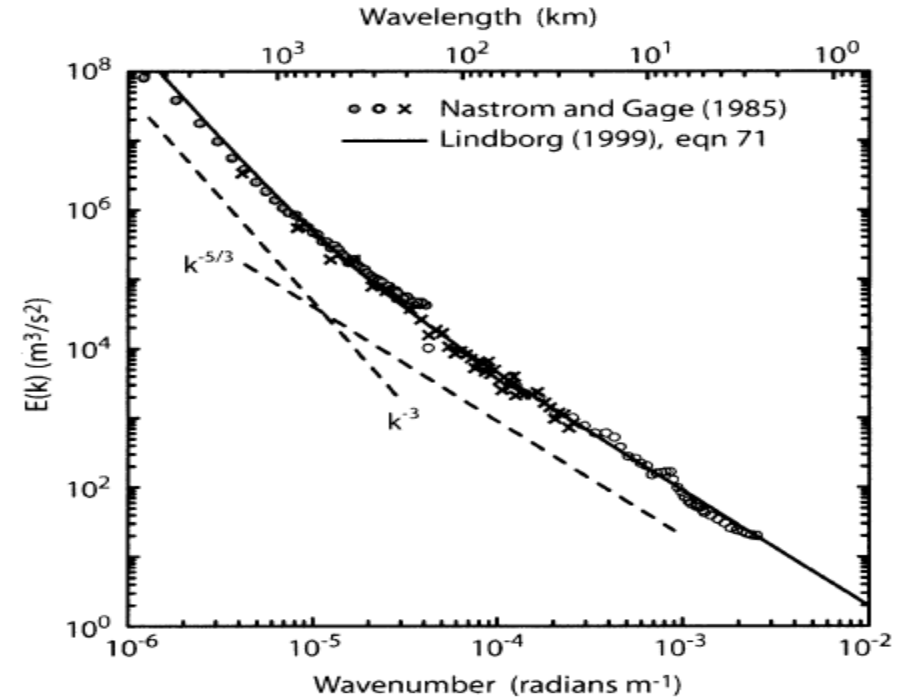


FIG. 1. Nastrom and Gage (1985) spectrum derived from the GASP aircraft observations (symbols) and the Lindborg (1999) functional fit to the MOZAIC aircraft observations.

- In order to fulfil the observed energy cascade from resolved scales to unresolved scales in GCMs, a **dissipation term** is added:

$$Dissip(\psi) = \frac{(-1)^{q+1}}{\tau} \nabla^{2q} \psi$$

Lateral dissipation in GCMs as a tool to pin the energy cascade

Figures from *Numerical Techniques for Global Atmospheric Models*, Lauritzen et al. (eds), 2010

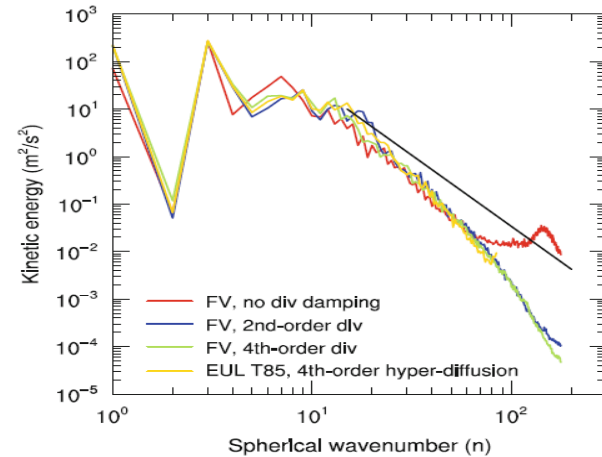
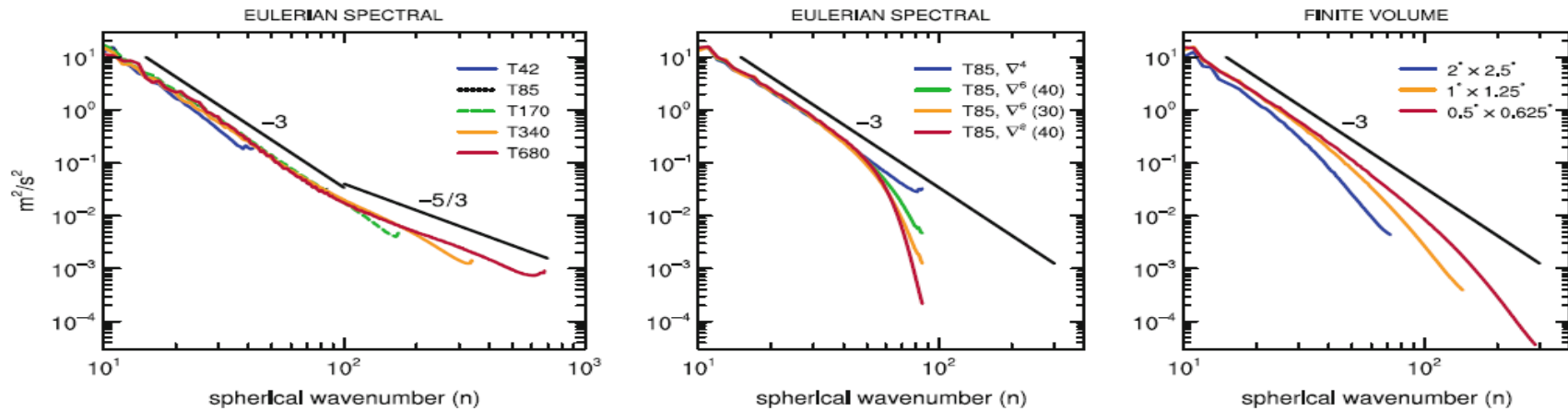


Fig. 13.4 250 hPa kinetic energy spectra as a function of the spherical wavenumber (n) in aquaplanet simulations from (left) CAM Eulerian spectral dynamical core with ∇^4 diffusion for different resolutions, (center) T85L26 Eulerian spectral dynamical with ∇^4 , ∇^6 and ∇^8 diffusion, and (right) CAM Finite Volume (FV) dynamical core for different $lat \times lon$ resolutions in degrees and 26 levels

Illustrative example of dissipation

- Simple 1D diffusion equation toy model:

$$\frac{dA}{dt} = \nu \frac{\partial^2 A}{\partial x^2}$$

- Von Neumann (Fourier mode) analysis

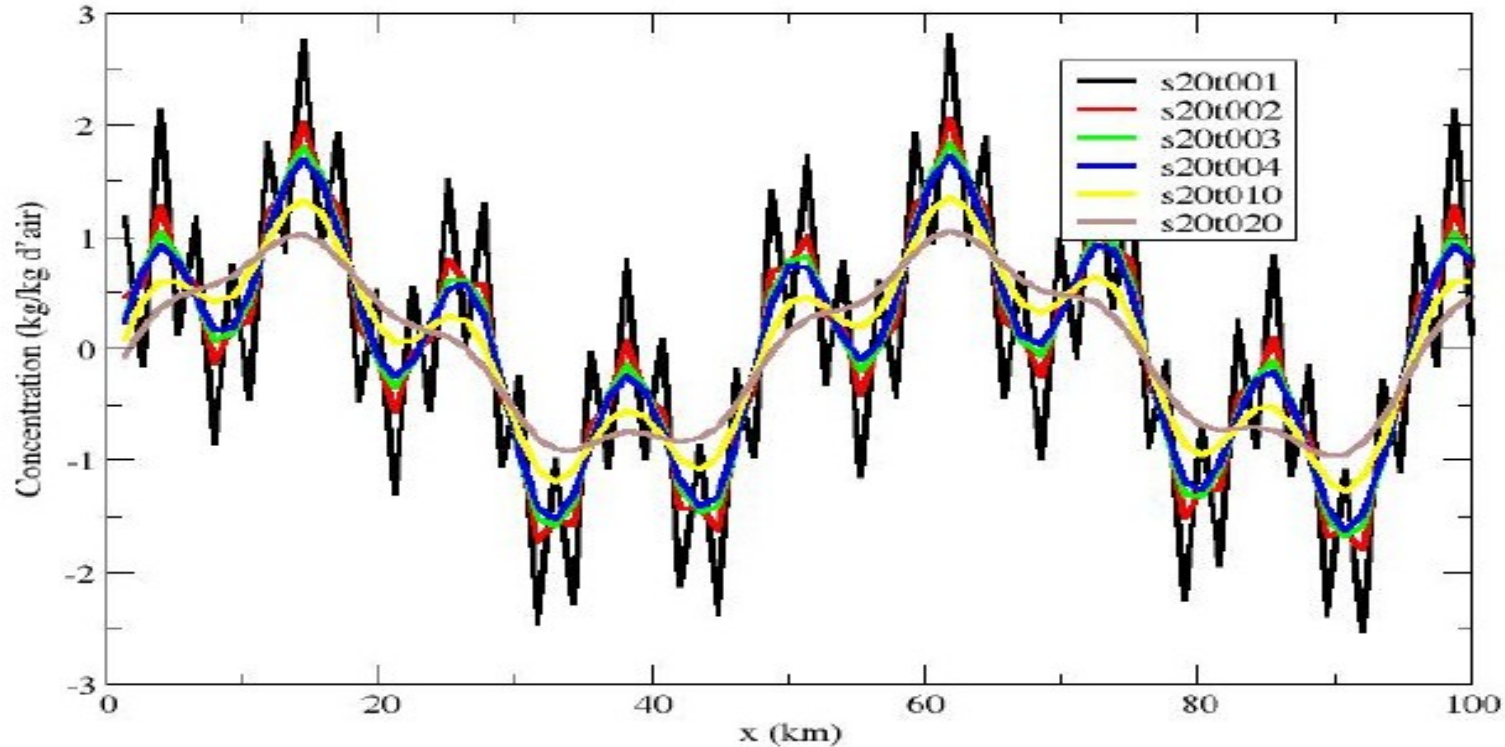
$$A_k(t) = a_k(t) \cdot \sin(kx)$$

- **Explicit Euler** time marching (with stability condition!):

$$a_k^{n+1} = (1 - \nu k^2 \Delta t) a_k^n$$

Note that mode damping is stronger for large k

Illustrative example of dissipation



Temporal evolution, from an initial condition consisting of 2 sine modes and an extreme (2 grid points wavelength) “numerical mode”

Controlling dissipation in LMDZ

- Parameters in file gcm.def:

dissip_period: Apply dissipation every dissip_period dynamical steps (or specify 0 to let model pick an appropriate value)

nitergdiv: number of iterations on velocity dissipation operator grad.div

nitergrot: number of iterations on velocity dissipation operator grad.rot

niterh: number of iterations on temperature dissipation operator div.grad

Usual values: nitergdiv=1, nitergrot=2, niterh=2

tetagdiv: dissipation time scale (s) for smallest wavelength for u,v (grad.div component)

tetagrot: dissipation time scale (s) for smallest wavelength for u,v (grad.rot component)

tetatemp: dissipation time scale (s) for smallest wavelength for potential temperature (div.grad)

values depend on horizontal resolution

Controlling dissipation in LMDZ

- Parameters in file gcm.def:

tetagdiv: dissipation time scale (s) for smallest wavelength for u,v (grad.div component)

tetagrot: dissipation time scale (s) for smallest wavelength for u,v (grad.rot component)

tetatemp: dissipation time scale (s) for smallest wavelength for potential temperature (div.grad)

optimal tetatemp values depend on horizontal resolution

- Moreover there is a multiplicative factor for the dissipation coefficient, which increases with model levels (see dyn3d_common/**inidissip.F90**), which can be controlled by flag “vert_prof_dissip”

The sponge layer

- In addition to lateral dissipation, it is necessary to damp vertically propagating waves (non-physically reflected downward from model top).
- The **sponge layer is limited to topmost layers** (usually 4) and added during the dissipation step.
- Sponge modes and parameters (gcm.def):

iflag_top_bound: 0 for no sponge, 1 for sponge over 4 topmost layers, 2 for sponge from top to 100 times topmost layer pressure

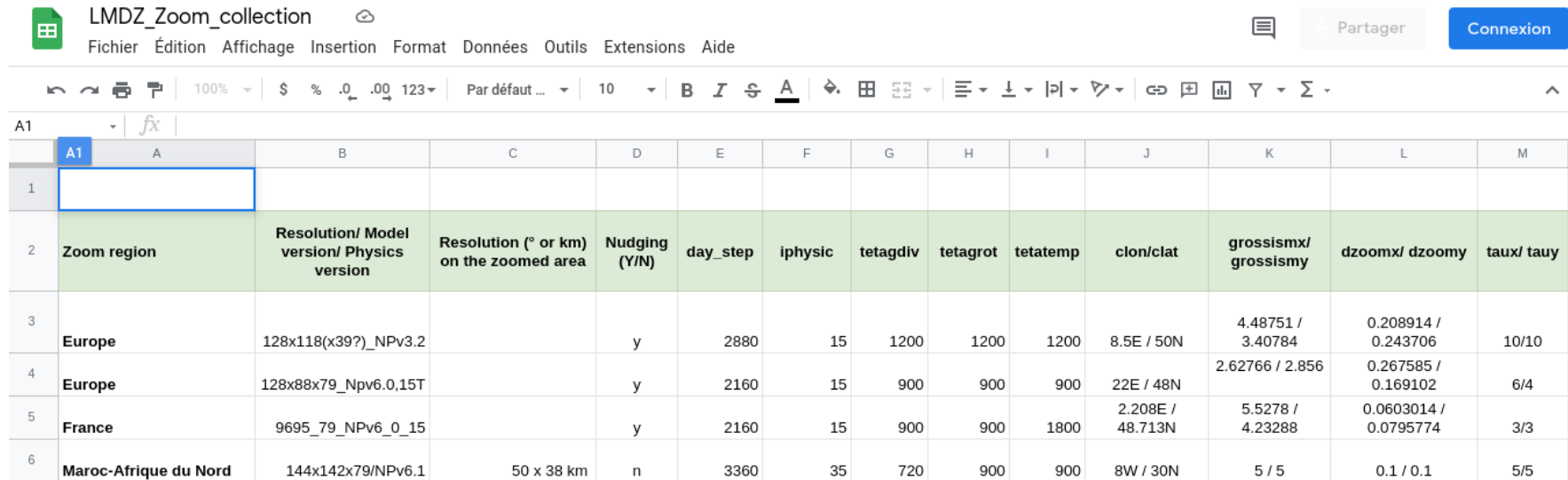
mode_top_bound: 0 for no relaxation, 1 to relax u,v to zero, 2 to relax u,v to their zonal mean, 3 to relax u,v and potential temperature to their zonal mean.

tau_top_bound: inverse of characteristic time scale at the topmost layer (halved at each successive descending layer)

Where to get typical values of the run.def/gcm.def parameters

- Check out the various examples of `gcm.def_*` files located in the **LMDZ/DefLists** subdirectory
- Some examples of specific cases of Zoomed simulation setups are detailed in a collaborative document; get the link by searching “Zoom collection” on LMDZPedia :

<https://lmdz-forge.lmd.jussieu.fr/mediawiki/LMDZPedia>



The screenshot shows a spreadsheet application with the title "LMDZ_Zoom_collection". The spreadsheet contains a table with 14 columns (A-M) and 6 rows. The first row (A1) is empty. The second row (A2) contains the following headers: "Zoom region", "Resolution/ Model version/ Physics version", "Resolution (° or km) on the zoomed area", "Nudging (Y/N)", "day_step", "iphysic", "tetagdiv", "tetagrot", "tetatemp", "clon/clat", "grossismx/ grossismy", "dzoomx/ dzoomy", and "taux/ tauy". The subsequent rows (A3-A6) provide data for "Europe", "France", and "Maroc-Afrique du Nord".

	A	B	C	D	E	F	G	H	I	J	K	L	M
1													
2	Zoom region	Resolution/ Model version/ Physics version	Resolution (° or km) on the zoomed area	Nudging (Y/N)	day_step	iphysic	tetagdiv	tetagrot	tetatemp	clon/clat	grossismx/ grossismy	dzoomx/ dzoomy	taux/ tauy
3	Europe	128x118(x39?)_NPv3.2		y	2880	15	1200	1200	1200	8.5E / 50N	4.48751 / 3.40784	0.208914 / 0.243706	10/10
4	Europe	128x88x79_Npv6.0,15T		y	2160	15	900	900	900	22E / 48N	2.62766 / 2.856	0.267585 / 0.169102	6/4
5	France	9695_79_NPv6_0_15		y	2160	15	900	900	1800	2.208E / 48.713N	5.5278 / 4.23288	0.0603014 / 0.0795774	3/3
6	Maroc-Afrique du Nord	144x142x79/NPv6.1	50 x 38 km	n	3360	35	720	900	900	8W / 30N	5 / 5	0.1 / 0.1	5/5

Rules of thumb for run.def parameters

- Time steps in LMDZ:

dynamical time steps: $dtvr = \text{daysec} / \text{day_step}$

physics time step: $dtphys = \text{iphysiq} * dtvr$

dissipation time step: $dtdiss = \text{dissip_period} * dtvr$

tracer advection time step: $dtvrtrac = \text{iapp_trac} * dtvr$

- Constraints to be aware of:

$dtvr$ limited by CFL for waves: $C_{\text{max}}.dt < \min(dx, dy)$

$dtrtrac$ limited by advection CFL: $U_{\text{max}}.dt < \min(dx, dy)$

$iphysiq$, $dtvrtrac$, $dissip_period$ **should be multiples of** $iperiod$

Rules of thumb for run.def parameters

- Constraints to be aware of (continued):
dissipation time step should be much smaller than dissipation timescales:
 $\text{dtdiss} \ll \text{tetatgdiv}, \text{tetagrot}, \text{tetatemp}$
- Changing time step with resolution on a regular grid:
 $\text{day_step}(\max(\text{iim}, \text{jjm})=N) \sim \text{day_step}(\max(\text{iim}, \text{jjm})=M) * M/N$
- Time step for a zoomed simulation, compared to regular grid:
 $\text{day_step}(\text{zoom}) \sim \text{day_step}(\text{regular}) * \max(\text{grossismx}, \text{grossismy})$