

Ongoing breakthroughs in convective parameterization

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Abstract Purpose of review: While the increase of computer power mobilizes a part of the atmospheric modelling community towards models with explicit convection or based on machine learning, we review the part of the literature dedicated to convective parameterization development for large-scale forecast and climate models.

Recent findings: Many developments are underway to overcome endemic limitations of traditional convective parameterizations, either in unified or multi-object frameworks : scale-aware and stochastic approaches, new prognostic equations or representations of new components such as cold pools. Understanding their impact on the emergent properties of a model remains challenging, due to subsequent tuning of parameters and the limited understanding given by traditional metrics.

Summary: Further effort still needs to be dedicated to the representation of the life cycle of convective systems, in particular their mesoscale organization and associated cloud cover. The development of more process-oriented metrics based on new observations is also needed to help quantify model improvement and better understand the mechanisms of climate change.

Keywords convective parameterizations for large-scale models · stochastic approaches · convective memory · mesoscale circulation · cold pools · process-oriented metrics

1 Introduction

Atmospheric moist convection results from the radiative cooling of the atmosphere, surface fluxes, forcing by large-scale motions, and the buoyancy associated with water phase changes. It happens at various spatial and temporal scales, from shallow convective cells to individual deep convective cells of a few kilometers lasting a few hours, to mesoscale convective systems of hundreds of kilometers lasting a few days, up to synoptic clusters and

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convection organized at the planetary scale such as the Madden-Julian Oscillation (MJO) or the Intertropical Convergence Zone (ITCZ), and is associated with the formation of liquid and ice hydrometeors. This makes convection particularly challenging to be represented in global models with a horizontal resolution ranging from a dozen to a few hundreds of kilometers, in which convection and its associated clouds are entirely or partly sub-grid scale. The collective effect of sub-grid cumulus clouds on the resolved large-scale variables has to be taken into account via a set of equations called a convective parameterization. Arakawa (2004) defines it as “an attempt to formulate the statistical effects of cumulus convection without predicting each individual cloud”. It is supposed to be valid over the various conditions encountered all over the globe but also to respond appropriately to large-scale modifications of the state of the atmosphere (Raymond, 1994). The convective parameterization is key in weather and climate simulations as it provides source terms for the dynamic equations of heat, humidity and momentum as well as macro- and microphysical properties of clouds for radiation. It determines the spatio-temporal distribution of precipitation and extreme weather events such as tropical cyclones, lightning or severe thunderstorms. It may also provide vertical profiles of mass flux and vertical velocity, as well as aerosol activation, chemical reactions and wet removal processes within clouds for the transport and evolution of hydrometeors, tracers, aerosols and chemical species.

The majority of the parameterizations currently at work in operational weather and climate models are updated versions of schemes built in the 1980s and 1990s (Arakawa and Schubert, 1974; Tiedtke, 1989; Gregory and Rowntree, 1990; Kain and Fritsch, 1990; Emanuel, 1991; Del Genio and Yao, 1993; Donner, 1993; Grell, 1993; Zhang and McFarlane, 1995). They are based on the mass-flux approach, which aims to explicitly represent the underlying convective physical processes, by decomposing the atmospheric column into different components (updrafts, downdrafts, subsiding environment). In recent years, however, the validity of traditional hypotheses has been questioned and there has been a resurgence of new original approaches to parameterization, with the use of stochastic components, revisiting of the mass-flux framework, or the introduction of new processes (cold pools, mesoscale circulations). Part of those developments have been motivated by the emerging question of scale-adaptability when model resolution enters the “gray zone” (5 to 25 km) and convective transport needs to be partitioned into parameterized and resolved components. It is noteworthy that some of these new developments have already been implemented in models, indicating a speeding of the model development process.

In the meantime, even if convective parameterizations are improved at the process scale, major model large-scale biases still remain in relation with clouds and precipitation, both in the mean state (e.g., the double ITCZ artifact, warm biases of mid-latitude continents and the eastern tropical oceans), and in variability from the diurnal to inter-annual time-scales (monsoon, MJO, ENSO, ...). Furthermore, much less attention has been paid to the clouds that accompany moist convection, even though shallow cumulus clouds represent the largest source of uncertainty in cloud feedback in the current generation of global climate models (Klein et al, 2017) and the anvils associated with deep convective cloud systems are a largely unconstrained contributor to climate change besides their well-documented increase in height with warming (Zelinka and Hartmann, 2011). The persistence of endemic model biases, combined with the inevitable trend toward higher model resolution, has led part of the community to abandon traditional convective parameterization development in favor of new approaches made possible by the increase of computing power, based on models resolving

convection and clouds at the kilometer scale (cloud resolving models; CRMs). The super-parameterization or multi-scale modeling framework (Randall et al, 2003; Tao and Moncrieff, 2009) is a hybrid approach that replaces a conventional convective parameterization by a 2D CRM, or even a 3D large-eddy simulation (LES) model (Grabowski, 2016) in each atmospheric column. Ultimately, the goal of scale-aware parameterization efforts is to eventually give way to global cloud resolving models that operate on scales that do not require any cumulus parameterization at all (Miyamoto et al, 2013). More recently, first attempts to develop a novel class of convective parameterizations based on machine learning, using a deep neural network trained by explicit simulations, have been tried with some skill to reproduce convective tendencies and clouds (Schneider et al, 2017; Gentine et al, 2018; O’Gorman and Dwyer, 2018). The question thus arises: Is it beneficial to continue to develop conventional cumulus parameterizations? In this paper based on the recent literature, we argue that long-term community effort should still be dedicated to traditional convective parameterization development for large-scale models for several reasons. First, overcoming the difficulties of traditional parameterizations has led to original and exciting developments in the last few years, demonstrating that various limitations of convective parameterizations can be overcome. Second, parameterizations summarize our understanding of physical processes and their interactions with the large-scale flow. For climate models in particular, this distillation of knowledge must be central to any strategy to understand climate change. Third, parameterizations potentially permit us to disentangle the role different physical processes play in weather and climate. In addition to this, there will still be decades before routine climate simulations can be run at convection permitting resolution.

It is often stated that a good climate model should be a good weather model, and indeed, short term error growth on weather time scales can reveal important parameterization deficiencies that lead to mean state biases (e.g. Van Weverberg et al, 2018). It has not been demonstrated, though, that a good weather model is sufficient to produce a reliable climate change projection. The goal of parameterization development is not just to improve model performance and get models to agree with each other, but also to understand why models predict the climate changes that they predict, estimate what confidence we can have in them, and articulate the reasons for that confidence (or lack thereof). To achieve this, we need to understand better how convection interacts with other processes to modulate the weather or climate and to quantify better the impact of related tuning parameters. It is also important to precisely define what a climate and weather forecast model should be able to simulate and find new ways to confront large-scale models with observations and high-resolution simulations that reveal the fidelity with which general circulation models (GCMs) represent the convective scale processes that they seek to parameterize.

In the following, we address three major challenges the convective parameterization development community is facing: 1/improve the representation of convective cloud ensembles 2/improve the representation of convective memory and organization, 3/improve the representation of convection to large-scale interactions.

2 Improve the representation of convective cloud ensembles

2.1 Unified versus multi-object frameworks

Atmospheric convection organizes at various vertical scales, going from boundary-layer depth to the full troposphere depth for deep moist convection as sketched in Fig 1. Boundary-layer convection is active all year long over the tropical oceans, where it produces cumulus and strato-cumulus clouds and is a key of the representation of the diurnal cycle over continents. The importance of the continuity between dry and cloudy boundary-layer convection has been recognized for decades (LeMone and Pennell, 1976), as well as the need to account for the non-local asymmetrical vertical transport between buoyant plumes and compensatory environment to allow the systematic up-gradient transport of potential temperature. Cumulus clouds are thus associated with the positive wing of a positively skewed distribution of sub-grid water vapor. Two kinds of approaches developed in the last decades which fulfill those constraints have started to be used in weather forecast and climate models: Combination of eddy diffusion with mass-flux schemes (often called EDMF) (Hourdin et al, 2002; Köhler et al, 2011; Hourdin et al, 2013; Riette and Lac, 2016; Bhattacharya et al, 2018) and third order turbulent schemes (Guo et al, 2014, 2015).

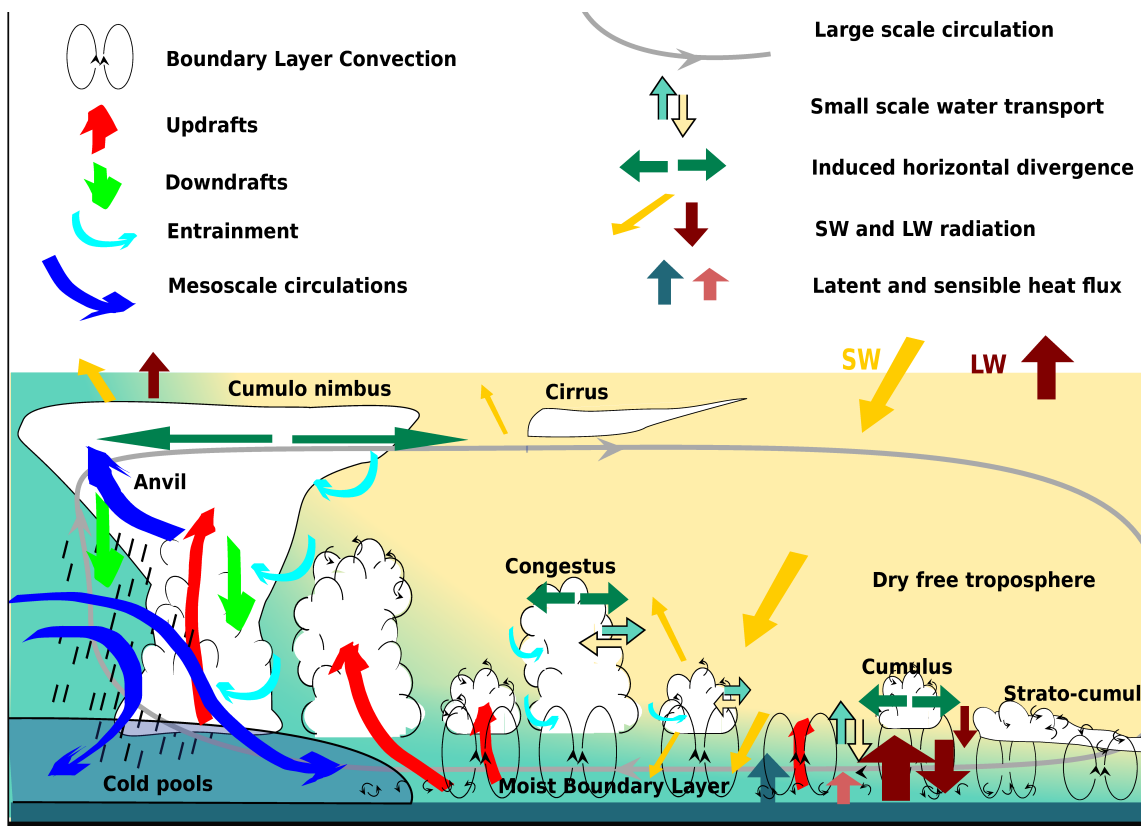


Fig. 1 Sketch of the various aspects of atmospheric convection and its interaction with radiation, humidity, surface fluxes and large-scale circulation which are targeted by parameterization developments, noticing that most of the aspects can co-exist within the same horizontal grid cell for a global climate model.

In regions where deep convection is active, observations show a trimodal distribution of clouds corresponding to (shallow) cumulus, (deep) cumulonimbus and the intermediate scale of congestus clouds (Johnson et al, 1999).

The various modes interplay in the Tropics. While shallow convection generally moistens and cools the free troposphere by detrainment of water vapor and cloud water that evaporates (Nitta and Esbensen, 1974), deep convection warms and dries the free troposphere by compensating subsidence and precipitation (Yanai et al, 1973). These three modes sketched in Fig 1 often coexist within domains that correspond to a model grid cell. Most models use mass-flux parameterizations for deep convective regimes, but with a variety of approaches and underlying conceptual models. Some modelers promote unified parameterizations for convection arguing that the physics of the convection is unique, and that a well-designed unified scheme should help represent the coexistence of various regimes (Gu er emy, 2011; Park, 2014a,b; D’Andrea et al, 2014) and work at grid resolutions where cumulus convection is permitted but not fully resolved (Arakawa and Wu, 2013; Wu and Arakawa, 2014; Kwon and Hong, 2017; Ong et al, 2017; Zhao et al, 2018). Others underline the contrasting nature of shallow and deep convection, in terms of organization and main drivers (surface fluxes for boundary-layer convection and condensation, rain evaporation, and sometimes wind shear for deep convection) and advocate for the coexistence of different mass-flux schemes, considering the coupling between those two objects as a parameterization itself (Rio et al, 2013). Representing cold pools as entities distinct from the deep convection scheme permits the introduction of a sub-grid partitioning of the otherwise homogeneous environment: the exterior of cold pools in which convective cells initiate, and the cold pool area in which convective rain falls and evaporates (Grandpeix and Lafore, 2010).

Note that there are also tentative works to extend parameterizations derived more directly from perturbation development of the basic equations (Guo et al, 2015; Storer et al, 2015). Whether they are called unified or not, it seems important that physics packages clearly target the various aspects and scales of convection, either by coupling different parameterizations, or including in a unique model the ingredients needed to represent the various conditions properly, including the coexistence of different regimes within the same grid box.

2.2 Size distribution of convective clouds

In practice, most schemes used in large-scale models rely on two updraft classes, one for shallow and one for deep convection, although several fully spectral ensemble approaches exist (Arakawa and Schubert, 1974; Donner et al, 2011). A triggering criterion, usually based on instability, moisture convergence or boundary-layer vertical velocity, is used to decide what parts of the scheme should be activated or not. Then closure relationships are used to relate convection intensity to large-scale variables. In recent years, several studies have targeted the representation of the co-existence of more numerous cloud types within a model grid cell with scale-adaptativity.

One approach consists in using a Markov chain to derive fractional areas covered by 3 cloud types (congestus, deep, stratiform) given probabilities of transition dependent on chosen large-scale variables (for example CAPE and mid-tropospheric dryness) as proposed by Khouider et al (2010). These fractions can be combined with vertical profiles of heating and moistening inferred from observations to compute convective tendencies (Deng et al, 2015; Goswami et al, 2017a,b). Others consider more numerous cloud types with transitions conditioned on mid-tropospheric vertical velocity and relative humidity (Peters et al, 2013) or on the vertical velocity averaged over the lower part of the troposphere (Dorrestijn et al, 2013). Both approaches are tested in GCMs by using

the deep convective area fraction to compute the closure of distinct versions of the Tiedtke (1989) mass flux scheme assuming a cloud-base vertical velocity of 1 m s^{-1} (Dorrestijn et al, 2016; Peters et al, 2017).

Rather than considering a finite number of cloud types, another approach derives a stochastic closure from an exponential PDF of mass-flux per cloud (Plant and Craig, 2008). The inferred probability of initiating a plume of a given radius is tested against random numbers, the ensemble-mean mass-flux being derived using a CAPE closure computed from large-scale variables with space-time averaging. The stochastic scheme has been coupled to the Kain-Fritsch (Keane and Plant, 2012; Keane et al, 2016) and the Zhang-McFarlane (Wang and Zhang, 2016) schemes. A similar approach has been derived for shallow convection and coupled to an EDMF scheme by Sakradzija et al (2016). Those approaches consider that clouds are randomly distributed over the domain, with no convective aggregation. This aspect has been taken into account by Hagos et al (2018) via another stochastic framework, not implemented in a GCM yet, which aims to predict the evolution of the size distribution of convective cells from a probability of growth, a probability of decay and an imposed relationship between the cloud base mass-flux and the convective area fraction.

Similar approaches are also used to modify the triggering of deep convection. A positive deep convective fraction provided by the multcloud model can be used as a triggering criteria (Peters et al, 2017). Rochetin et al (2014a) compute a distribution of sizes and velocities of thermal plumes from a mean updraft issued from an EDMF scheme. Deep convection is triggered if the strongest thermal has a sufficient vertical velocity to overcome the convective inhibition and if one thermal within the grid cell is larger than a given threshold. Deep convection is initiated randomly according to the specified distribution of thermals. The stochastic components used in those frameworks in principle allow for scale-adaptability (Keane et al, 2014; Sakradzija et al, 2016).

2.3 Vertical structure of the updraft

Whatever the approach used for triggering and closure and the number of cloud types, there is still the need to compute the vertical profile of cloud properties. Using directly observed profiles of heating and moistening as in Goswami et al (2017b) may improve the coupling of convection with the large-scale dynamics but it has some limitations for deriving mass-fluxes for vertical transport of chemical species and aerosols or for simulating the evolution of the vertical profiles of convective tendencies under climate change. In the other approaches cited above, classical formulations of entrainment and detrainment are applied to each cloud type. This is convenient as in practice it remains challenging to find an exact same formulation of entrainment that is valid for both shallow and deep convection (Del Genio and Wu, 2010; Zhang et al, 2016). Direct computation of entrainment in LES models suggests that determining entrainment rates using the bulk approximation underestimates effective entrainment by a factor of 2 (Romps, 2010). The presence of shells around convective clouds (Heus and Jonker, 2008; Glenn and Krueger, 2014) implies that the air entrained at the edges of cumulus clouds does not have the properties of the mean environmental air as commonly assumed in parameterizations, but rather is a mixture of cloudy and environmental air (Zhang et al, 2016). However lower entrainment rates used in bulk-plume schemes might be an efficient way to simplify the more complex mixing occurring at cloud edges, the important aspect being to take into account properly the dilution of the updrafts by entrainment (Hannah, 2017). For example,

Becker et al (2018), using a CRM, find that when convection aggregates, the entrainment rate actually increases in the lower troposphere due to the enhancement of turbulence, but that the dilution of cloud air decreases because the updrafts entrain the more humid air that surrounds aggregated updrafts.

While some use deterministic entrainment applied to several plumes with different cloud-base mass-flux, other studies have shown that models considering entrainment as a stochastic process capture more accurately the variability of cloud properties seen in LES (Romps and Kuang, 2010; Böing et al, 2014) than models using multiple updrafts with deterministic entrainment rates. This may be achieved by traditional episodic mixing and buoyancy sorting schemes (Raymond and Blyth, 1986; Emanuel, 1991; Grandpeix et al, 2004), provided the vertical profile of mixture distributions and the properties of updraft and environmental air that mix can be constrained from CRM or LES. Alternatively, Romps (2016) proposes a new approach tested in a single column model (SCM) based on an updraft initialized with a single set of properties at the first model level and made of convective parcels entraining stochastically as a Poisson process. An attempt to couple an earlier version of this scheme (Romps and Kuang, 2010) in an EDMF scheme has been made in a 1D framework (Sušelj et al, 2013).

Another issue concerns the vertical profile of vertical velocity, which is required for coupling with microphysics and eventually incorporating chemistry and aerosol effects on convection (Lee et al, 2009; Donner et al, 2016). While some schemes do not solve any equation for the vertical velocity, most of the ones doing so rely on slightly different implementations of the formulation proposed by Simpson and Wiggert (1969). However this equation does not take into account properly the effect of pressure perturbations, that influence maximum velocity and cloud top height, an issue addressed recently theoretically by Peters (2016); Morrison (2016a,b). Also, the role of buoyant and inertial acceleration in the triggering of convective mass-flux by cold pools has been addressed by Jeevanjee and Romps (2015).

The various cloud types that co-exist within a model grid cell are also involved in the life cycle of convective systems, an additional issue in that case being to represent the self-sustaining behavior of convection.

3 Improve the representation of convective memory and organization

3.1 Departure from quasi-equilibrium and memory

Even though several observational or CRM studies have shown the limitations of the quasi-equilibrium assumption between convection and the large-scale forcing introduced by Arakawa and Schubert (1974) (Zhang, 2002, 2003; Donner and Phillips, 2003), most climate model convective parameterizations still rely on it. However quasi-equilibrium breaks down for domains smaller than 250 km and when the forcing period becomes less than 30h for organized convection (Jones and Randall, 2011) and 12h for more scattered convection (Davies et al, 2013). In addition, even if quasi-equilibrium is observed over large domains or periods of time, it might best be used as an emergent behavior to evaluate a closure assumption rather than being a closure itself (Del Genio and Yao, 1993; Arakawa, 2004). The quasi-equilibrium hypothesis also does not account for convective memory effect (Davies et al, 2009). A way to include this effect is to introduce prognostic variables in the convection scheme. Pan and Randall (1998) propose a prognostic cumulus kinetic energy used to compute the cloud-base mass-flux. Some approaches introduce prognostic variables whose evolution depends on rain evaporation to

modify the entrainment or closure formulations (Piriou et al, 2007; Hohenegger and Bretherton, 2011; Mapes and Neale, 2011; Chen and Mapes, 2018), without representing explicitly the physical processes responsible for these feedbacks (downdrafts, cold pools). Recently, Colin et al (2018) show using idealized CRM experiments that convective memory is mostly carried out by low-level thermodynamic structures and is enhanced when convection is organized at the mesoscale.

In recent years, attempts have been made to represent more explicitly the complex physics of processes involved in the self-sustaining behavior of convection. Grandpeix and Lafore (2010) introduce a full parameterization of an ensemble of cold pools whose thermodynamic properties and fractional coverage are prognostic and driven by unsaturated downdrafts of the Emanuel (1991) deep convection scheme. In turn, the cold pool parameterization provides a lifting energy and power used for the triggering and closure computations of the deep convection scheme. The number of cold pools per unit area is shown to be a key parameter of the parameterization (Grandpeix et al, 2010) which is different over land and ocean. Via a parameterization of intermediate complexity, Del Genio et al (2015) regulate the occurrence of weakly entraining convection by cold pools whose evolution is computed via the introduction of two prognostic variables (the cold pool area and pressure depth) and by relaxing their thermodynamic properties to those of the undisturbed boundary-layer on different time scales over land and ocean. In a unified framework, Park (2014a) introduces memory via a prognostic treatment of the sub-grid cold pool and mesoscale organized flow and their feedback on convective updrafts.

Without taking into account cold pools, some of the stochastic schemes presented in section 2.2 also introduce some memory effect via the use of Markov chains or a master equation (Peters et al, 2013; Hagos et al, 2018). An extended version of an EDMF scheme with a prognostic treatment of plume area fractions has also been proposed and tested in a 1D model by Tan et al (2018).

3.2 Towards mesoscale organization

Despite their importance for sustaining convection, only one operational GCM has represented any aspect of mesoscale convective systems (MCSs) (Donner et al, 2011), although several ideas for parameterizing various MCS effects have been proposed (Alexander and Cotton, 1998; Mapes and Neale, 2011; Khouider and Moncrieff, 2015; Yano and Moncrieff, 2016, 2018; Moncrieff et al, 2017). Doing so requires identifying which process controls the organization of deep convective cells: Wind shear, a moist free troposphere, and a long-lived moist boundary layer seem to promote “sustainability” of convection that precedes organization (Yuter and Houze, 1998; Houze, 2004; Schumacher and Houze, 2006). Traditionally, wind shear has received the most emphasis (Rotunno et al, 1988; Moncrieff and Liu, 1999), but convection can organize in weak shear as well (Houston and Wilhelmson, 2011). Cold pools generated by convective downdrafts are often a first step in deepening and sustaining convection (Khairoutdinov and Randall, 2006; Böing et al, 2012; Del Genio et al, 2012; Schlemmer and Hohenegger, 2014) and create memory as well. However, they do not explain organization in all cases; in the presence of a very moist boundary layer and low convective inhibition, gravity waves are probably the relevant organizing/propagation mechanism (Huang, 1990; Lane and Moncrieff, 2015; Grant et al, 2018). This has not been accounted for yet in any GCM.

Cold pools or gravity waves can be a precursor but actual organization implies subsequent development of a mesoscale circulation in the stratiform rain region (Houze, 2004). Sustaining convection to produce the thick anvil that generates the stratiform rain region places a premium on simulating the interaction between convective dynamics and the microphysics of convective ice, which determines detrainment and can affect climate sensitivity (Zhao et al, 2016). Radiative heating in the resulting anvil initiates a mesoscale updraft that nucleates more ice and ultimately drives a mesoscale downdraft via melting and evaporation of rain below (Houze, 2004). Cloud-resolving model simulations suggest that the magnitude of the updraft and downdraft can be constrained by the diabatic heating they produce (Del Genio et al, 2012).

Climate GCMs still operate at horizontal resolutions (50-200 km) at which MCSs are primarily deterministic and there is no horizontal propagation of convective systems, unless the resolved flow creates it. NWP models however run routinely at fine resolutions at which stochastic behavior is important. One suggestion for incorporating organization into a stochastic framework in such models is via cellular automata (Bengtsson et al, 2013).

3.3 Associated clouds and precipitation

Precipitation and clouds are involved in both the spatial organization and the time evolution of convective systems. Most microphysics schemes used in convection schemes are based on simple formulations that compute rain rates from liquid and ice thresholds and precipitation efficiency depending on pressure or updraft temperature involving tuning parameters (Zhang and McFarlane, 1995; Emanuel and Zivkovic-Rothman, 1999; Mauritsen and Stevens, 2015). Recently a two-moment diagnostic parameterization that calculates mass mixing ratio and number concentration of four hydrometeors (cloud liquid water, cloud ice, rain, snow) based on Morrison and Gettelman (2008) has been coupled to convection schemes (Song et al, 2012; Storer et al, 2015; Storer et al, 2015). Schemes including an equation for the vertical velocity can implicitly include the effects of microphysics in updrafts by specifying particle size distributions (PSDs) and size-fall speed relationships to partition precipitation and detrainment (e.g. Del Genio et al, 2005). When applied with PSDs and fall speeds derived from recent field experiments (Mitchell et al, 2011; Heymsfield et al, 2013), this approach can produce tropical anvil ice water paths in good agreement with satellite observations (Elsaesser et al, 2017). Realistic vertical velocities are also needed to adequately control aerosol activation of both liquid droplets and ice crystals (Feingold, 2003; McFiggans et al, 2006; Kay and Wood, 2008).

Regarding the computation of cloud cover, even if some early diagnostic cloud schemes have let deep convective clouds interact with radiation (Slingo, 1987), the lack of a proper coupling between convective suspended and falling hydrometeors and radiation is still responsible for important radiative biases of climate models (Li et al, 2016). Recently, several attempts have been made to improve convective cloud radiative effects based on LES analysis that have shown that the double-gaussian PDF is the most accurate to represent the sub-grid scale cloud structure for both shallow and deep convection (Bogenschutz et al, 2010; Perraud et al, 2011) if the required input moments can be predicted or diagnosed accurately. Gaussian (Qin et al, 2018) and bi-gaussian (Jam et al, 2013; Hourdin et al, 2013) PDFs have been implemented in some models to represent

shallow convective clouds with the PDF variances diagnosed from the turbulent and shallow convective processes. Double-Gaussian PDF are also used by Storer et al (2015) to extend the assumed probability density function method to deep convection. There is not yet a satisfactory way to treat the cloud cover associated with the anvil clouds formed by detrainment, even though prognostic stratiform cloud schemes that account for it exist (Tiedtke, 1993; Del Genio et al, 1996; Tompkins, 2002).

The most controversial aspect of convective clouds is their sensitivity to changes in aerosol concentration. In polluted environments, updrafts should nucleate more and thus yield smaller droplets, which would suppress rain by reducing fall speeds and increasing rain evaporation. However if the cloud deepens and the smaller droplets are more easily lifted into the mixed-phase region, they may increase ice formation and precipitation and invigorate convection instead (Rosenfeld et al, 2008). Fan et al (2018) suggest that in cleaner environments, condensation onto ultrafine aerosols and the resulting latent heat release is the invigoration mechanism. Khain (2009) summarizes possible aerosol and environmental influences on convective clouds, classifying them into regimes based on how condensate generation by aerosols competes with precipitation loss, to explain seemingly conflicting previous conclusions about aerosol suppression vs. invigoration. Unfortunately, it has been difficult to determine whether aerosol effects on convection are important, because in models, uncertainties in parameterized cloud microphysics greatly exceed the aerosol effect (White et al, 2017), while in real world case studies, meteorological and thermodynamic effects overwhelm any aerosol signal (Varble, 2018).

Disentangling the respective role of convective processes in the emergence of large-scale features proves to be difficult. Both radiative effects of convective clouds and convective transport play a role in the interactions between convection and the large-scale flow. In the last section, we discuss how the successive phases of the convective life-cycle interact with the large-scale and relevant observations to evaluate models at the process level.

4 Improve the representation of convection to large-scale interactions

4.1 Shallow convection

Shallow convective clouds are the saturated part of thermals initiated at the surface (LeMone and Pennell, 1976). Combining a mass-flux approach with a diffusion scheme within the boundary-layer produces a more efficient vertical transport in dry and shallow convective regimes. The associated enhanced export of the evaporated surface water to the dry troposphere and enhanced compensating subsidence enhances surface drying, in better agreement with observations over land (Cheruy et al, 2013; Diallo et al, 2017). Over ocean, this leads to an enhancement of the drying of the boundary-layer in subsiding regions. This increases the surface evaporative cooling, the underestimation of which was shown to explain part of the East Tropical Ocean (ETO) warm biases in CMIP simulations (Hourdin et al, 2015b). Shallow convection is also important for the injection to the free troposphere of trace species emitted at the surface (Locatelli et al, 2015), which are then further transported by deep convection (Folkens et al, 2006; Donner et al, 2007), as well as for coupling with chemistry (Nie et al, 2016). The mass-flux transport of horizontal momentum improves the representation of the diurnal cycle of

near-surface winds, from desert areas, with a strong effect on dust lifting (Hourdin et al, 2015a), up to the Antarctic Plateau (Vignon et al, 2018).

The meridional distribution of the (shallow or deep) cloud radiative forcing has been shown from analysis of CMIP (Xiang et al, 2017) and idealized aquaplanet experiments (Dixit et al, 2018; Talib et al, 2018) to control the latitudinal distribution of tropical rainfall. This may be key to explain the systematic tendency of global models to overestimate convective rainfall south of the equator over the Atlantic and Pacific oceans (the so-called double ITCZ issue), as shown for example by Qin and Lin (2018). The underestimation of shadowing by strato-cumulus also in part explains systematic warm biases over ETO in CMIP models (Hourdin et al, 2015b).

Shallow convective clouds moreover explain a significant fraction of the range of cloud feedback estimates which dominate the spread in Equilibrium Climate Sensitivity (ECS) to greenhouse gases (Bony and Dufresne, 2005; Zelinka et al, 2016; Geoffroy et al, 2017; Vial et al, 2018). Aside from, or coupled with radiative effects, the literature increasingly emphasizes the role of vertical water vapor transport by convection (Sherwood et al, 2014; Klein et al, 2017; Cesana et al, 2018) or the condensate transport that determines detrainment of ice (Zhao et al, 2016) in the control of the ECS. The feedback from these clouds appears to be positive, due to drying of the subsiding regions in which these clouds occur. Observed inter-annual variations support such a conclusion (Klein et al, 2017). LES models subjected to an imposed climate change, though, find the feedback to be only weakly positive (Zhang et al, 2013).

4.2 Transition from shallow to deep convection

The transition from shallow to deep convection has received considerable attention, especially in connection with the representation of the diurnal cycle of precipitation (Couvreur et al, 2015) and medium-range predictability associated with the MJO (Klingaman et al, 2015). Various approaches have led to postpone diurnal deep convection initiation over land rather by modifying the convective entrainment formulation (Stratton and Stirling, 2012), modifying the CAPE closure (Bechtold et al, 2014) or introducing a separate EDMF scheme used to trigger deep convection (Rio et al, 2009; Rio et al, 2013). While models are known to rain too frequently (Dai, 2006) and to overestimate the autocorrelation with previous day precipitation over land (Roehrig et al, 2013), stochastic triggering is able to increase the day-to-day variability of precipitation over land (Rochetin et al, 2014b). The triggering criterion also strongly influences the MJO (Peters et al, 2017). An aspect that deserves more attention is the representation of the effect of sea breeze in deep convection initiation over islands, a key mechanism for the representation of precipitation over the maritime continent (Birch et al, 2015).

Once deep convection is initiated, a key issue is to be able to represent the correct partitioning between the various convective regimes and their respective effects. The often neglected congestus phase can influence the representation of spatio-temporal precipitation variations (Hirota et al, 2014). Numerous data analyses document the transition from shallow to deep convection and its association with the sensitivity of convection to free tropospheric humidity via entrainment (Morita et al, 2006; Benedict and Randall, 2007; Holloway and Neelin, 2009; Riley et al, 2011; Del Genio et al, 2012; Kuo et al, 2018). Still unresolved, though, is the tendency

for models with strong intraseasonal variability caused by this sensitivity to also have strong positive tropical rain biases (Kim et al, 2011), suggesting that GCMs do not yet strike the right balance between suppressed and vigorous deep convection. Also, several studies point for the consequences of the misrepresentation of convective transport or microphysics on model biases: impact of momentum transport on large-scale fields (Zhang and Cho, 1991; Wu et al, 2007; Lane and Moncrieff, 2010; Orr et al, 2010; Woelfle et al, 2018), of convection induced gustiness on the structure of the ITCZ (Harrop et al, 2018), of convective entrainment on the double ITCZ syndrome (Oueslati and Bellon, 2015), of the assumed temperature at which convective updraft supercooled liquid water glaciates on the Southern Ocean shortwave radiation bias (Kay et al, 2016).

While recent developments have been directly based on in-situ data (Dorrestijn et al, 2013; Peters et al, 2013), surprisingly little effort has been made to validate the properties of either the spectrum of convective cells produced by parameterizations, or the “bulk plume” that represents their collective effects. Information about convective vertical velocities and/or mass fluxes now exists (Collis et al, 2013; Giangrande et al, 2013; Kumar et al, 2015; Giangrande et al, 2016; Masunaga and Luo, 2016; Labbouz et al, 2018). Combined with particle size distributions in convective outflow from field experiments and ice water path from satellites, detrainment can be evaluated (Elsaesser et al, 2017). Satellite observations now constrain the areal coverage and spatial distribution of convective cells (Fu et al, 1990; Schumacher and Houze, 2003; Sassen and Wang, 2008), and top height distribution (Takahashi et al, 2017), which might provide useful indirect constraints on entrainment.

4.3 Transition from deep to organized convection

Mesoscale convective systems (MCSs) are responsible for many important weather and climate phenomena: (1) Top-heavy heating that influences the tropical general circulation (Schumacher et al, 2004); (2) The maintenance and propagation of convection, which controls its diurnal peak (Nesbitt and Zipser, 2003); (3) Extreme precipitation events (Doswell et al, 1996; Houze and Churchill, 1987; Mathon et al, 2002; Moseley et al, 2016); (4) Coupling with large-scale tropical wave phenomena such as the Madden-Julian Oscillation (Xu and Rutledge, 2015) and possibly African easterly waves (Ventrice and Thorncroft, 2013); (5) The sign of convective momentum transport (Moncrieff et al, 2017); (6) Tropical cloudiness and energy balance (Kiehl, 1994; Rossow et al, 2005; Hartmann et al, 2018); (7) The cloud height feedback on climate (Hartmann and Larson, 2002).

The physical mechanisms responsible for the interaction of convection with these phenomena vary. In some situations, cold pools and their associated gust fronts may determine the organization and propagation speed of MCSs (e.g. Rotunno et al, 1988). In others, the system moves at a speed different from that of the cold pools (Lane and Moncrieff, 2015). Planetary-scale propagating disturbances usually have MCSs embedded in their disturbed phase. For some, such as the MJO, the MCSs are considered integral to the existence of the larger-scale phenomenon, since the mesoscale updrafts of the MCSs produce an extensive anvil cloud shield whose longwave heating is thought to be the ultimate source of moist static energy for the MJO (e.g. Andersen and Kuang, 2012; Kim et al, 2015). The mesoscale updraft-downdraft combination in the stratiform rain regions of many MCSs produces a dipole heating anomaly that shifts the overall heating profile of the system upwards, intensifying the upper-level large scale tropical circulation (Schumacher et al, 2004). GCMs currently represent

all these interactions via their cellular convection or grid-scale condensation schemes since they do not represent MCSs.

The question for any convective organization scheme is the extent to which it reproduces the real-world behavior of MCSs: the extent to which their existence and behavior is deterministic vs. stochastic; the convective-stratiform partitioning of rain; the shape of the heating profile; and the clouds they produce. Validation of cold pool parameterizations has relied mostly on anecdotal information (Rozbicki et al, 1999), but statistics of observed (Zuidema et al, 2017; Schiro and Neelin, 2018) and simulated (Feng et al, 2015) downdrafts and cold pools now exist. Databases of satellite-derived MCS lifecycles and tracks also exist (Fiolleau and Roca, 2013) as do estimates of their macrophysical and microphysical properties (Bouniol et al, 2016), and of the convective-stratiform partitioning and heating profiles (Lang and Tao, 2018; Shige, 2009). These have not been widely utilized by the parameterization community. To the extent that such products have been used in model intercomparisons (e.g. Dai, 2006), model-data agreement must be interpreted as the result of compensations for missing physics in other parts of the parameterization since most models do not represent the mesoscale processes that actually determine convective-stratiform partitioning and the shape of the heating profile.

5 Discussion and conclusions

While a part of the community doubts parameterization will ever work for convection, recent developments of new concepts and approaches synthesized here rather draw the picture of ongoing breakthroughs in convective parameterizations in recent years. Consequently, while the computationally intensive approaches will continue to progress toward the ultimate goal of reducing the problem to parameterizations of microphysics, turbulence, and radiation, we anticipate vigorous parallel development of traditional parameterizations to include all the processes that need to be represented to understand the mechanisms of climate change. Indeed, several shortcomings of traditional parameterizations are being tackled by ongoing developments in different groups: the underestimation of vertical transport by shallow cumulus clouds by the development of unified approaches for boundary-layer dry and moist convection; the hypothesis that cumulus clouds cover a small fraction of a grid cell by scale-aware approaches; the mean bulk hypothesis by the consideration of a spectrum of cloud types of different sizes or stochastic entrainment; the quasi-equilibrium hypothesis by the introduction of new prognostic variables to allow both quasi-equilibrium and departures from it, depending on the situation; the too-short lived convective systems by the development of approaches to represent cold pools or, even still few, mesoscale circulations. The improvement of the representation of clouds associated with convection, including more fundamental approaches to precipitation efficiency and microphysics, certainly deserves to be reinforced, since it is as much involved in the interaction between convection and the large-scale as convective transport and is essential for any attempt to simulate aerosol-convection interactions. The role a convection scheme has to play at gray-zone resolutions also needs to be clarified, as some studies highlight that models without any cumulus parameterization perform better than models with parameterized convection at those resolutions, despite the fact that gray zone models do not explicitly resolve convective cells. This is verified for example when comparing models run at horizontal resolutions from 4 to 40 km with either parameterized or explicit convection over West

Africa (Birch et al, 2014) and over the Indian ocean (Holloway et al, 2012, 2013, 2015). A fundamental issue to be addressed is how cumulus parameterizations deteriorate feedbacks relevant for large-scale features such as the West-African monsoon or the Madden Julian Oscillation and what key behaviors cumulus (and other) parameterizations need to have in order to add value at these resolutions.

It will certainly still take time until developments concerning all those aspects of convection are developed and implemented in every operational model. Indeed, side by side with this enthusiastic picture, there is still the feeling that progress is slow in climate models. Several issues may deserve more attention. Only a few recent studies explore in a systematic way the sensitivity of convective parameterizations to the value of the very uncertain free parameters (Bernstein and Neelin, 2016). The use of “meta-models” or “emulators”, developed in the community of Uncertainty Quantification (Sacks et al, 1989; Williamson et al, 2015) may help explore more efficiently the parameter space and address an issue which is often informally discussed among climate modeling groups: the fact that significant changes in the rainfall distribution after changes of convective parameterization are often overruled by the subsequent re-tuning of the clouds’ radiative effect (Hourdin et al, 2017). This is a fundamental result of the fact that in steady state, atmospheric latent heating must balance radiative cooling, an unavoidable consequence of global coupled modeling (Schmidt et al, 2017). One relevant example is the apparent importance of radiative forcing for the presence or absence of features such as the double ITCZ bias. One consequence of this coupling between latent heating and radiative cooling is that if non-convective processes are not realistically represented, convection will have to be tuned in unrealistic directions to compensate. This may partly explain why potential improvements to cumulus parameterizations identified in studies that look only at some aspect of the convection itself often do not find their way into operational models. To make progress, a more integrated view of convection within the larger scope of other processes that affect the hydrologic cycle and energy cycle is required. For example, Webb et al (2015) show that removing the cumulus parameterization from a sample of GCMs significantly changes the cloud feedback in some of them, but not the overall spread in cloud feedback, and that the changes in feedback are greatly reduced when the models clouds’ are rescaled to the observed cloud radiative effect in the current climate.

The weight given to parameterization development vs. intercomparison exercises deserves serious discussion by the community as well. CMIP exercises have become central in climate modeling, and they are indeed the only way to properly measure the progress in fully integrated climate models and evaluate the component of uncertainties in climate modeling that comes from model “structural errors”. Many studies, some of them discussed in Section 4, have made full use of CMIP simulations (and in particular of pairs of atmosphere-only and coupled simulations (Xiang et al, 2017; Hourdin et al, 2015b) which constitute the essence of the CMIP6 “DECK”) to identify important processes which should be better parameterized. Comparatively, studies that try alternative solutions (such as those proposed in Sections 2 and 3) in real climate models are fewer. When such studies test changes in convection, taking into account a proper retuning of the model (Zhao et al, 2016), they underline the difficulty in disentangling the coupling between the various aspects of the physics parameterizations (convective transport, cloud micro- and macro-physics, convective moisture transport). The use of a hierarchy of model configurations is probably a promising path to follow: SCM configurations for detailed comparison with LES, 3D nudged simulations to get synoptic observations in phase with in situ observations, Transpose-

AMIP approaches to look at the fast response of model errors and finally atmosphere-alone and coupled models. Analysis of the recent literature suggests that those approaches are more efficient when used simultaneously in a single model to disentangle processes, than across models as a basis for new intercomparison exercises, the human cost of which is probably underestimated.

Another issue is to get rid of the tyranny of traditional metrics in favor of process-oriented metrics (Jakob, 2010). Large-scale phenomena such as the ITCZ, MJO, etc., as well as mean state biases, are often used to validate cumulus parameterizations, but emergent behavior of this kind can only be considered a necessary, rather than sufficient, condition for having confidence in any cumulus parameterization. The use of traditional metrics has not led to narrowing the range of model predictions of climate change, suggesting that these metrics do not target behaviors of the models that matter for their intended use. Even worse, reliance on such metrics discourages experimentation with novel approaches, which when implemented in models often make the models worse before they are fully understood. Likewise, “emergent constraints,” though intended to bypass the problems of traditional mean state bias metrics, often have their own problems. Caldwell et al (2018) find that most proposed constraints on equilibrium climate sensitivity do not pass one or more tests for plausibility. Before large-scale emergent behavior is evaluated, cumulus parameterization performance needs to be measured against observational constraints closer to the process level. As discussed in section 4, the convective life cycle, which can be described as a progression from shallow to deep to organized mesoscale convection (Mapes et al, 2006), provides a useful framework for thinking about cumulus parameterization at the process level and for connecting parameterization development to explicit models and observations.

The last issue is the scheme complexity. The recent developments described above often result in an increasing complexity of the schemes with some counterparts: a difficulty to fully describe one particular scheme in the literature, making difficult the discussion of concepts and comparison of approaches, the increased probability of conceptual errors and bugs in codes that nobody else other than the developer can identify, and finally the difficulty to make new approaches work in real weather prediction or climate models (due to cost, numerical instability, difficulty in correcting undesirable behaviors when coupled to the dynamics). Reaching real progress in the representation of convection in climate models is probably not only a question of new brilliant ideas or concepts, but also a multi-dimensional issue that concerns physics understanding of convective systems, observations, tuning, methodological aspects, numerical issues and coding as well as human organization of modeling teams and priorities of research programs. Nonetheless, the progress that has occurred in recent years, despite these roadblocks, suggests that the field of cumulus parameterization development is more dynamic and promising than it seemed to be some decades ago.

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